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Persistent Supercooled Drizzle at Temperatures below -25°C Observed at
McMurdo Station, Antarctica

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Abstract

The rarity of reports in the literature of brief and spatially limited observations of drizzle at temperatures below -20°C suggest that riming and other temperature-dependent cloud microphysical processes such as heterogeneous ice nucleation and ice crystal depositional growth prevent drizzle persistence in cold environments. In this study, we report on a persistent drizzle event observed by ground-based remote-sensing measurements at McMurdo Station, Antarctica. The temperatures in the drizzle-producing cloud were below -25°C and the drizzle persisted for a period exceeding 7.5 hours. Using ground-based, satellite, and reanalysis data we conclude that drizzle was likely present in parts of a widespread cloud field, which stretched more than ~ 1000 km along the Ross Ice Shelf coast. Parameter space sensitivity tests using two-moment bulk microphysics in large-eddy simulations constrained by the observations suggest that activated ice freezing nuclei (IFN) and accumulation-mode aerosol number concentrations aloft during this persistent drizzle period were likely on the order of 0.2 L^{-1} and 20 cm^{-3} , respectively. In such constrained simulations, the drizzle moisture flux through cloud base exceeds that of ice. The simulations also indicate that drizzle can lead to the formation of multiple peaks in cloud water content profiles. This study suggests that persistent drizzle at these low temperatures may be common at the low aerosol concentrations typical of the Antarctic and Southern Ocean atmospheres.

Key Points

- * Drizzle was observed over McMurdo at cloud $T < -25^{\circ}\text{C}$ for more than 7.5 hours in the presence of ice nucleation and growth
- * Atmospheric conditions necessary for the observed drizzle are evaluated using an LES model with two-moment microphysics
- * Drizzle at cloud $T < -25^{\circ}\text{C}$ may be common at the low aerosol concentrations typical of the Antarctic and the Southern Ocean atmospheres

1. Introduction

Drizzle at sub-freezing temperatures has been detected in various studies, both in mid-latitudes (e.g., Cober & Isaac, 2012; Ikeda et al., 2007; Rosenfeld et al., 2013) and the polar regions (e.g., Lawson et al., 2001; Verlinde et al., 2013). In convective clouds, it is possible that drizzle and raindrops may be lofted from the 0°C isotherm up to temperatures as low as -20°C when the mean updrafts are > 6 m/s (Carey & Rutledge, 2000; Zipser & Lutz, 1994). However, in stratiform clouds or convective clouds with weaker updrafts, the formation and persistence of drizzle may depend on the condensate supply (e.g., by adiabatic lifting) and collision-coalescence rate, C_i , exceeding the sum of heterogeneous freezing, ice crystal depositional growth, and/or riming rate, D_i (e.g., Cober et al., 1996; Ikeda et al., 2007; Lawson et al., 2001; Rosenfeld et al., 2013), that is, $C_i > D_i$.

Collision-coalescence of cloud droplets is the predominant mode of drizzle production (e.g., Rasmussen et al., 2002), generally requiring low cloud condensation nuclei (CCN) concentrations and/or high liquid water path (LWP) for the development of sufficiently broad droplet size distributions (e.g., Comstock et al., 2004; Pruppacher & Klett, 1997, pp. 581–598). The formation of drizzle without an active collision-coalescence process has also been reported when droplet number concentrations are extremely low (Mauritsen et al., 2011). Such number concentration requirements make this process more likely to be active in a maritime environment (e.g., Cober et al., 1996, 2001), or potentially, in stable environments where updrafts are very weak.

The $C_i > D_i$ requirement suggests that the ice crystal concentrations should also be sufficiently low (hence, smaller D_i) to enable drizzle formation and persistence (e.g., Rasmussen et al., 2002). Nonetheless, drizzle can be maintained at greater D_i if C_i is greater still. For example, while Cober et al. (1996) measured ice crystal number concentrations of 0.08 L^{-1} in supercooled drizzle regions with weak updrafts (< 0.1 m/s), Ikeda et al. (2007) reported ice crystal number concentrations < 2 L^{-1} in drizzle formation regions, which experienced consistent updrafts stronger than 0.5-1 m/s.

Ice crystal concentrations sufficiently low for drizzle formation can be found at temperature ranges outside the secondary ice formation region (Geresdi et al., 2005), either in a relatively pristine environment with low ice freezing nuclei (IFN) concentrations (e.g., Antarctica; Bromwich et al., 2012), or in cloud in which IFN were depleted from the drizzle formation regions (from a Lagrangian perspective) by ice precipitation (Ikeda et al., 2007; Rosenfeld et

al., 2013). However, the nucleation efficiency of IFN sharply increases with decreasing temperature and/or increasing ice supersaturation (Meyers et al., 1992; Phillips et al., 2008), which steeply rises at water saturation with decreasing temperatures. Therefore, drizzle formation is expected to be rarer and of shorter duration in colder clouds. For example, Rosenfeld et al. (2013) detected drizzle at -21°C along a flight track of 53 km, whereas Lawson et al. (2001) reported a smaller 10 km drizzle pocket (~ 100 m below cloud-top) at an even lower temperature of -25°C , which was completely depleted at the same location only two minutes later (and 30 m lower). The lowest temperature at which drizzle has been reported, to the best of our knowledge, is -28°C , observed by Korolev et al. (2002) over a 3 minute period (18 km flight track).

In this study we document the occurrence of persistent drizzle formation for over 7.5 hours in a shallow, mixed-phase cloud layer with temperatures ranging between -25°C and -29°C (from cloud base to cloud top), which was observed using remote-sensing measurements at McMurdo Station, Antarctica (77.85°S , 166.72°E ; ~ 70 m above mean sea level). The measurements were gathered as part of the U.S. Department of Energy (DOE) Atmospheric Radiation Measurement (ARM) Facility's (Mather & Voyles, 2013) 2016 ARM West Antarctic Radiation Experiment (AWARE) field campaign (Lubin et al., 2015; Witze, 2016). In addition, we use large-eddy simulations (LES) constrained by these measurements to evaluate the conditions required for drizzle formation at these temperatures.

2. Drizzle Event Observations

The drizzle was detected at McMurdo Station on August 16, 2016, between 05:05 and 12:45 UTC (a total of 460 min, all at night), in a persistent mixed-phase cloud that lasted for nearly 39 hours at temperatures ranging between -23°C and -29°C , which was first observed ~ 10 h prior to the period discussed in this study. A deep cyclone centered over the Southern Ocean to the north of McMurdo Station throughout the event generated northeasterly winds across the eastern Ross Sea, which advected warm and moist marine air towards the Ross Ice Shelf (RIS; **Figure 1a**). The cyclone-driven wind confluence with the RIS southwesterly katabatic flow resulted in convergence along the RIS coast, including around McMurdo Station (relatively weaker winds in **Figure 1a**; see Figure S1 in the Supporting Information for high-resolution flow patterns around Ross Island). Widespread enhanced top of atmosphere (TOA) upwelling longwave (LW) radiation footprints were measured in this coastal region during this event by the Clouds and the Earth's Radiant Energy System (CERES; dashed red shape in **Figure 1b**) instrument on NASA's Aqua satellite.

Figure 2a-b show the High Spectral Resolution Lidar (HSRL; Eloranta, 2005) particulate backscatter cross-section (β_p) and linear depolarization ratio (LDR) measurements during the drizzling period. The liquid-bearing cloud regions are indicated by high β_p and low LDR (predominantly $> 5 \cdot 10^{-5} \text{ m}^{-1} \text{ sr}^{-1}$ and < 0.08 , respectively). Note that the HSRL backscattered signal is often extinct below the mixed-phase cloud layer top (marked by purple dots), which is determined by a Ka-band ARM Zenith Radar (KAZR) moderate sensitivity (MD) mode signal-to-noise ratio (SNR) threshold of -16 dB (following the methodology in Silber et al., 2018a). The ceilometer cloud base height (Morris, 2016) and HSRL liquid cloud base height (LCBH; see Silber et al., 2018b, 2018c) data products, which are depicted by white and black dots, respectively, mostly display little variability near the mixed-phase layer base. The agreement between these two data products is good, with a typical difference of ~ 50 - 100 m (cf. Silber et al., 2018c). We refer hereafter to the HSRL data product when discussing the LCBH, as this cloud base height algorithm is more sensitive (relative to the ceilometer data product) to the level of condensation in this optically-thin (per unit volume) mixed-phase cloud layer.

Concurrent moderate β_p ($\sim 10^{-5.3}$ - $10^{-4.5} \text{ m}^{-1} \text{ sr}^{-1}$) and low LDR (< 0.2) values are frequently observed down to ~ 300 m below the LCBH (**Figure 2a-b**; see also **Figure 3a-b**), with higher LDR values closer to the surface. These higher LDR values indicate precipitating ice. The low LDR values just below the LCBH suggest that the backscattered lidar signal from these air volumes is dominated by nearly spherical scatterers, for example, falling drizzle drops (diameter < 0.5 mm per Pruppacher & Beard, 1970). While specular reflection by plate ice crystals may induce similar low LDR values (see Appendix A in Silber et al., 2018c), the cloud layer temperatures (T ; **Figure 3h**, left panel) were well below the plate crystal growth regime (Bailey & Hallett, 2009). Aerosol particles and cloud droplets may also produce low LDR, but the measured β_p values are too high and low for aerosol particles and cloud droplets, respectively (e.g., Silber et al., 2018a, Figure 2). Moreover, the temporal variability of these β_p -LDR features is significantly patchier than that of a typical liquid-bearing cloud layer. Therefore, we deduce that these features are indicative of drizzle drops, similar to the remote-sensing signatures observed by Verlinde et al. (2013).

To further support our conclusion, we examine measurements from a single hour (10-11 UTC; **Figure 3**), coinciding with a radiosonde release (the only in-situ measurements obtained during this drizzle event). During this hour (and a large fraction of the whole drizzle event; see **Figure 2c**), the KAZR reflectivity factor (**Figure 3c**) shows frequent fall streaks

below the LCBH. These regions of higher reflectivity (e.g., 10:30 UTC and 10:52-10:55 UTC) correspond with an increase in LDR with decreasing altitude right below LCBH (**Figure 3b**). The concurrent reflectivity and LDR patterns suggest that these fall streaks are composed of ice particles that nucleate near (or at) cloud top (**Figure 3c**), which have an adequate surface area to hasten evaporation of the low LDR liquid hydrometeor population. The most upward Doppler velocity in the KAZR Doppler velocity spectrum with signal above the noise floor, V_{dmax} , is depicted in **Figure 3d**. V_{dmax} can be used as an estimation of the vertical air motion in regions where there are droplets large enough (and/or at a sufficient concentration) to produce a detectable signal and small enough to have terminal velocities that are slow relative to the air motion. V_{dmax} may be upward-biased because of turbulent broadening, which for the data above the lowest LCBH in Figure 2d is equivalent to $\sim 0.2\text{--}0.4$ m/s (Appendix A). Suitable droplet populations can be found within mixed-phase cloud layers, although not necessarily at the LCBH where the droplet volume can be too small to produce backscattered signals above the KAZR noise floor when there is no drizzle. Verlinde et al. (2013, Figure 9) showed for a weakly drizzling case that V_{dmax} decreased slowly across the LCBH, with the evaporating small drizzle drops (diameter $< 200\ \mu\text{m}$) providing the backscatter necessary for detection by KAZR. In these drizzle-bearing air-volumes below the LCBH, the liquid water content (LWC) is significantly lower than within the cloud layer, and hence, the effect of spectral broadening biases described in Appendix A are negligible. When the total volume of drizzle drops in the radar volume decreases below the KAZR detectable signal, the V_{dmax} signal comes from the faster-falling ice in the volume, and hence, no longer serves as a proxy for the air motion. In regions with stronger liquid hydrometeor signals below the LCBH (**Figure 3a-b**) higher saturation ratios driven by weak updrafts of up to ~ 0.1 m/s (**Figure 3d**) may help counter the desiccation effects of the ice. During the stronger ice showers, when liquid hydrometeors are mostly absent, V_{dmax} no longer represents the vertical air motion. Rather, the observed negative values represent the Doppler velocity of the slowest falling ice particles.

The KAZR spectra (velocity bin resolution of 2.34 cm/s) show distinct signatures of a slow-falling hydrometeor population at and below the LCBH, both in the vertical spectra profiles (e.g., at $\sim 10:22$ UTC; circled in **Figure 3e**) and in the temporal evolution of the spectra at a fixed elevation close to the LCBH (e.g., at 10:00-10:30 UTC; **Figure 3f**, circled areas). This Doppler spectra population cannot be attributed to humidified aerosols (cf. van Diedenhoven et al., 2011), the reflectivity of which would not exceed the KAZR noise floor. The fall

velocity of the hydrometeors in the circled population is < 0.5 m/s, typical for drizzle drops with diameters < 150 μm (Lamb & Verlinde, 2011, Figure 9.5). The small decrease in the population's fall speed relative to V_{dmax} seen with decreasing altitude (e.g., at $\sim 10:22$ UTC; **Figure 3g**) is consistent with drops evaporating below the LCBH. The ice population peak fall velocity is approximately constant with decreasing height at these drizzle-bearing altitudes, implying that riming is likely not a very active process at these heights. Such a lack of a riming signature suggests insufficient LWC and/or number concentrations (e.g., Jensen & Harrington, 2015). Although individual ice crystals growing in the ice-supersaturated conditions beneath the LCBH should see a monotonic increase in fall speed along Lagrangian trajectories, such a signal is not readily evident in the instantaneous snapshot represented in the spectrograph of many such trajectories captured at various stages of evolution. A similar lack of an active riming process is indicated in the first 90 min of the drizzle event when maximum LWP was measured (**Figure 2f**).

The balloon sounding occurred during a period in which the HSRL drizzle signature was observed down to ~ 850 m. As shown by these measurements (**Figure 3h**), the mixed-phase cloud layer T ranged between -25°C and -29°C (left panel). Beneath the LCBH, the relative humidity with respect to water (RH_w ; middle panel) declines from saturation with decreasing height. The specific humidity (q) profile suggests that the mixed-phase cloud and the layer below it (down to ~ 1000 m) are not well mixed as not uncommonly expected in stratiform mixed-phase clouds (e.g., Fridlind et al., 2012), although the equivalent potential temperature (θ_e ; left panel) does show a tendency for low stability ($d\theta_e/dz \approx 0.33$ K km^{-1} on average in the altitude range of 1.0-1.6 km). Gaps in the lowest liquid cloud base intermittently reveal another cloud base that is ~ 100 m higher, suggesting that LWC does not exhibit a linear profile with a persistent vertical structure. Lack of vertical shear of the horizontal wind, as well as the weak wind speed (~ 2 m/s; right panel) also do not support consistently strong mixing.

Taken together, the observations discussed above, that is, the liquid hydrometeor desiccation by stronger ice showers below LCBH, the low Doppler fall velocities of an independent hydrometeor population, and the sub-saturation RH_w below the LCBH, support our conclusion that persistent drizzle formation occurred throughout this long-lasting event.

The earlier work on drizzle formation in mixed-phase clouds suggested that updraft strength and duration may play a role (e.g., Cober et al., 1996; Ikeda et al., 2007). Because McMurdo Station is situated close to complex orography, larger-scale forcing of the vertical velocity

field should be considered to understand all the factors forcing the formation of drizzle. Using the Doppler radar observations, we perform a wavelet analysis (Torrence & Compo, 1998) of V_{dmax} 100 m above the LCBH throughout the entire drizzle event, to examine whether the drizzling periods correspond with local wave activity. We assume that cloud droplets at this height within the mixed-phase layer are consistently large enough to be detected by the KAZR and small enough to represent an estimation of the vertical air motion. The bias in V_{dmax} identified in Appendix A is inherently removed in the wavelet analysis. The “main” liquid-bearing cloud base used as a reference for this analysis (depicted in **Figure 2c**) is identified in intermittent periods of multiple liquid-bearing layers (see **Figure 2a-b**, **Figure 3a-b**) by examining its height consistency over time. LCBHs (at times with multiple liquid-bearing layers) that depart from the median ± 2 standard deviations (SDs) of all LCBH values between the 45th and 55th percentiles in 20 min windows (with 10 min overlap) are rejected from the analysis. This filtering step, which produced LCBH consistency, is followed by the removal of all residual multi-layer LCBHs above the lowest layer; **Figure 2a-b** show that the “main” liquid-bearing layer is predominantly the lowest one in multi-layer cases. Finally, LCBH values are linearly interpolated at times where LCBHs are not detected or no HSRL data are available.

The wavelet analysis spectrogram (**Figure 2d**) shows statistically significant oscillations (colored patches) at different times and periods. Statistical significance is determined here by comparison of the wavelet spectrum with a red noise spectrum and utilizing a chi-squared test (Torrence & Compo, 1998) at the 99% level, following the methodology of Miles and Verlinde (2005). Note that significance tests are not performed outside the cone of influence (COI; solid black curve in **Figure 2d**), as the spectral power in that region is affected by zero padding of the data necessary for the analysis procedures (see Torrence & Compo, 1998).

The spectrogram is divided into three main bands, subjectively selected to represent from the smallest to largest time scale the cloud-scale circulation and inhomogeneity, gravity wave (GW) scale oscillations, and synoptic-scale vertical air motion. The band limits (dashed and solid blue lines in **Figure 2d**) are located at 17 s, a variable ~4-7 min, 60 min, and 8 h. The lowest period threshold was chosen to correspond with the second-highest frequency received by the wavelet analysis; hence, the fastest oscillating component, which includes the turbulent eddies, is overlooked in this analysis. The variable 4-7 min periodicity corresponds with the apparent Brunt-Vaisala frequency (BVF) at McMurdo Station under the assumption of a horizontal wavelength of 4 km. This apparent BVF was calculated 0-200 m above the

LCBH using the Antarctic Mesoscale Prediction System (AMPS, on a 1.1 km horizontal mesh; Powers et al., 2012) model output bi-linearly interpolated to McMurdo Station's coordinates. The horizontal wavelength value matches small-scale topographic features to the northwest-west and south of McMurdo Station, which may act as a source of GWs induced by orographic lifting, given the prevailing large-scale wind direction (and intensity over the McMurdo region (Figure S1 in the Supporting Information)).

The GW band (band 2) shows frequent statistically significant oscillations, more often than the two other bands. The reconstruction of the band time series (**Figure 2e**; cf. Torrence and Compo, 1998) indicates that these GW oscillations can reach amplitudes of up to ~0.5 m/s. However, the variable GW-driven air motion, in sign and strength, does not necessarily produce drizzle-favored conditions (e.g., continuous lifting), which were reported in previous supercooled drizzle case studies (e.g., Ikeda et al., 2007). This absent correlation is evident here by the fact that drizzle signatures are observed below the LCBH both during positive and negative oscillation periods (compare **Figure 2e** to **Figure 2a-b**), whether these oscillations are intense (e.g., 11:45-12:00 UTC) or weak (e.g., 08:30-09:30 UTC). Some influence of the GWs on the liquid water may be evident, as the cloud liquid water (**Figure 2f**) does seem to deplete during intense oscillations (e.g., 06:30-07:30 and 11:00-12:00 UTC). The main processes that govern the desiccation of the cloud water, that is, entrainment of drier air at cloud top, ice growth and sedimentation from the cloud-containing layer, in-cloud collisions or riming and subsequently faster droplet scavenging via drizzle and ice precipitation, cannot be straightforwardly determined. Nonetheless, stronger ice precipitation at the base of the mixed-phase cloud layer is observed during these periods, as evident in the high LDR and concurrent higher reflectivity in **Figure 2b-c**, respectively. Further investigation of GW impact on cloud layers is beyond the scope of this study.

The persistent supercooled drizzle formation at liquid-bearing cloud temperatures $< -25^{\circ}\text{C}$ is evident both during periods with high ($> 100 \text{ g/m}^2$) and very low ($< 25 \text{ g/m}^2$) LWP (**Figure 2f**). It is produced either with or without substantial ice nucleation and growth (e.g., 06:00-07:00 and 07:35-08:20 UTC in **Figure 2a-c**, respectively), and without a clear connection to GW activity (**Figure 2e**). As the synoptic-scale air motion (band 3 in **Figure 2e**) reaches amplitudes of up to 0.13 m/s, which are comparable to the air motion detected locally over McMurdo at the LCBH during periods of drizzle (**Figure 3d**), and assuming based on the flow patterns (**Figure 1a**) and a back-trajectory (BT) analysis (Figure S2 in the Supporting Information) that the aerosol particle concentrations in the air masses above McMurdo

Station and the RIS coast are similar, we conclude that it is plausible that drizzle formation was extensive over the region of enhanced TOA upwelling longwave along the RIS coast measured by CERES (dashed red shape in **Figure 1b**).

We note that widespread drizzle formation along the RIS coast is not consistent with ERA5 (Hersbach & Dee, 2016) analysis (Figure 1a). However, a comparison of the water vapor mixing ratio profile during the drizzle event between the radiosonde sounding measurements from McMurdo and the corresponding grid cell in the ERA5 reanalysis indicates that ERA5 underestimates water vapor over this region by an average of ~13% for altitudes between 0-2 km. This underestimation could lead to an underestimation of the horizontal and vertical extent of condensate or the generation of ice clouds instead of liquid-bearing clouds (see Silber et al., 2019). For example, the ERA5 LWP at McMurdo Station (right inset panel in **Figure 1a**) is ~7.5 times smaller than the value derived from the microwave radiometer (Morris, 2006) measurements at McMurdo during the same hour (10:00-11:00 UTC; **Figure 2f**). ERA5 does produce low amounts of liquid water along the RIS coast but does not maintain liquid water from one hour to the next at specific locations. Because there is less upwelling LW radiation at TOA in ERA5 along the RIS coast compared to CERES measurements (not shown), we argue that the enhanced upwelling radiation measured by CERES along the RIS coast indicates the occurrence of widespread, relatively tenuous, mixed-phase clouds. This interpretation is further supported by the space-borne Cloud-Aerosol Lidar with Orthogonal Polarization (CALIOP) measurements (Winker et al., 2003) during 10:48-10:50 UTC, which provide evidence that the enhanced TOA upwelling longwave radiation along the RIS coincides with low-level mixed-phase cloud layers (dashed red shape in **Figure 1b**; Figure S2 in the Supporting Information). The variability in mixed-phase cloud heights seen in the CALIOP measurements (Figure S2a,b) explains the patchiness in the enhanced upwelling radiation region. While the existence of such clouds does not necessarily mean that drizzle was produced in large sectors of this widespread cloudy region, it certainly is a possibility.

3. Drizzle Event Modeling

A leading question that arises from these observations is how low the activated IFN and aerosol particle concentrations need to be to allow persistent drizzle formation at these low temperatures. To provide estimates for these two parameters, we use the observations described above to constrain model simulations. Our approach is meant to represent a plausible, quasi-idealized scenario under which the observed cloud formed.

For large-eddy simulations, we use the Distributed Hydrodynamic Aerosol and Radiative Modeling Application (DHARMA) model (Stevens et al., 2002), which includes a dynamic Smagorinsky scheme for subgrid-scale mixing (Kirkpatrick et al., 2006). DHARMA performs comparably to other LES models in case studies of mixed-phase clouds using bulk or bin microphysics; here we use a bulk microphysics scheme that prognoses two moments of the size distributions for each of three hydrometeor species: cloud water, rain, and cloud ice. Under these conditions the rain is very light and mean raindrop sizes quite small, so we will refer to rain and drizzle interchangeably.

To avoid the complications of specifying large-scale horizontal flux divergences, we translate the LES domain with the mean wind in the cloud layer, which is elevated and decoupled from the surface (see **Figure 3h** and **Figure 6a**). Vertical air motions within a cloud layer correspond to some combination of large-scale ascent, GW oscillations, and turbulence, with implications for the relationships between aerosol particles, droplet number concentrations, and drizzle formation. Plausible scenarios of large-scale vertical motion are first obtained from BT analysis of ERA-Interim (Dee et al., 2011) reanalysis fields using the Hybrid Single-Particle Lagrangian Integrated Trajectory model (HYSPLIT; Stein et al., 2015) over the range of the observed cloud depth during the target observation time (10-11 UTC; see **Figure 3**). All BTs show a period of notable ascent during the 9 h window prior to the target observation time (**Figure 4a,b**). A series of simulations for window periods from 6 to 24 h revealed that results at 10 UTC converge for all simulations with BT windows exceeding 9 hours. The LES applies the evolving profile of large-scale vertical motions from the 9 h back trajectories as source terms for prognostic variables computed from first-order upwinding (cf. Ackerman et al., 2009) rather than including as part of the flow, which would implicate mass continuity issues.

To derive a plausible initial sounding for the LES, we apply the reversed mean BT ascent rates (i.e., an effective descent) in a DHARMA 1-D configuration for a 9 h run, initialized with the 10:41 UTC radiosonde T and q profiles, in some resemblance to the approach taken by Field et al. (2012). The observed profiles are smoothed with spline fits using a smoothing parameter of 10^{-6} , with fitting points selected to also remove the cloud signatures and produce a supersaturated layer (see the green curves in **Figure 3h**). The removal of the cloud signatures from the atmospheric profiles prior to the 1-D backward run allows us to examine, using the LES, the cloud formation in an elevated, stable, sub-saturated moisture layer via radiative and/or large-scale ascent driven cooling. This approach is motivated by the work of

Simpfendorfer et al. (2019), who showed that cloud processes themselves contribute to the characteristics of observed sounding. The objective here is to allow the LES to form and develop the cloud-induced mixed-layer, as well as the cloud-top temperature and moisture inversions. The backward 1-D run accounts only for condensational adjustments with respect to liquid water. Radiative heating is ignored. At the end of this 9 h “backward” run, no cloud water remains in the atmospheric profile, consistent with a “pre-cloud” environment. The resulting T and q profiles, along with uniform ~ 2 m/s horizontal winds are used to initialize the LES (data are given in Table S1 and illustrated in Figure S3 in the Supporting Information). The LES setup and physics options are described in more detail in Appendix B. It should be noted that the results presented below are not sensitive to halving the grid spacing, doubling the domain size, reducing the time step duration, nudging to a vertically uniform 0 m/s horizontal wind, or incorporating vertical shear of the horizontal wind, from 2 m/s at 1.5 km to 9 m/s at 2 km, matching the observations in **Figure 3h**.

Hourly evolution of an LES run for activated IFN number concentrations (N_{IFN}) of 0.2 L^{-1} and accumulation-mode aerosol particle concentrations (N_A) of 20 cm^{-3} is depicted in **Figure 5** (an animation file with 1 min temporal resolution is provided in the Supporting Information). Non-negligible amounts of cloud water mass mixing ratio (q_c) are first observed after 2.75 h (**Figure 5a**). The cloud-top height (CTH) increases by ~ 360 m from the cloud layer’s formation until the end of the simulation. About 100 m of this CTH increase is attributable to vertical displacement of the profile forced by the large-scale ascent, while the rest of the CTH increase is associated with cooling of the overlying air by adiabatic expansion and LW radiative cooling, and weak cloud deepening through entrainment (cf. Randall, 1984) after the cloud becomes turbulent (**Figure 5h**).

Ice and drizzle first appear 20 min and 1.5 h after the initial cloud formation, respectively (**Figure 5b,c**), and their downward transport via sedimentation acts to reduce the total water mass mixing ratio (q_t) from their formation region near cloud top (**Figure 5d**). However, while the ice reaches the surface as a result of the supersaturated relative humidity with respect to ice (RH_i) nearly down to the surface (**Figure 5e**), the drizzle evaporates below the LCBH. The net effect of the drizzle is to transport moisture from cloud top to just below the LCBH, resulting in a local maximum q_t underneath the cloud, and a negative q_t gradient over the depth of the cloud layer (**Figure 5d**).

At ~ 6.5 h, the radiative cooling near cloud-top is unable to overcome the stability there. However, because the cloud is semi-opaque in the LW, radiative cooling occurs throughout

the depth of the cloud layer (**Figure 5f**), thereby promoting enough instability in mid-cloud (**Figure 5g**) to initiate turbulent mixing (**Figure 5h**). This turbulent mixing, while growing in magnitude and vertically, acts to reduce the negative- q_t and positive- θ_e gradients, as evident in **Figure 5d,g**. This diminution of vertical gradients leads to greater q_c near cloud top (**Figure 5a**), at the top of the turbulent layer (**Figure 5h**), and a decrease at the base of the turbulent layer (**Figure 5a,h**), by transporting warm, moist air upwards and cool, dry air downwards. The higher q_c near cloud top increases the LW opacity there, which enhances cloud-top radiative cooling (**Figure 5f**). The enhanced radiative cooling, in turn, overcomes the stability in the upper layer of the cloud, allowing positive feedback to develop between stronger turbulence, increasing q_c at cloud top, and cloud-top radiative cooling (**Figure 5a,f,h**).

At the same time, the q_c reduction at the base of the turbulent layer just above the LCBH (**Figure 5h**) splits the cloud layer in the sense that a double-peaked cloud structure develops (**Figure 5a**). The double-peaked structure is maintained for nearly two hours until the end of the simulation, though it does not persist much beyond that in longer simulations (not shown). The upper q_c peak is supported by the ample cloud-top radiative cooling that invigorates the turbulence below, while the evaporation of the lower q_c peak is hindered by the local drizzle-driven q_t maximum (**Figure 5d**). We thus speculate that drizzle evaporating in quiescent air may play a key role in the redistribution of moisture necessary to support the formation and persistence of multi-peaked cloud water content profiles, which are occasionally observed both over Antarctica (e.g., Silber et al., 2018c, Figure 1) and the Arctic (e.g., Silber et al., 2018c, Figure S2; Verlinde et al., 2013, Figures 1 and 8).

After its first appearance in the baseline model run, drizzle is continuously formed (**Figure 5c**), regardless of the ascent rates (**Figure 4b**). Drizzle is also produced in additional LES sensitivity tests, in which the same N_{IFN} and N_A values are used while the ascent rates are set to 0 cm/s from different times of the model run, before and after the initial cloud formation, until the end of these 9 h runs (Figure S4 in the Supporting Information). However, the initial appearance of drizzle is delayed by up to several hours in these sensitivity tests (not shown). When the model is initialized with higher N_{IFN} and/or N_A values, drizzle production is inhibited (Figure S5), especially without continuous ascent (Figures S4). Therefore, we postulate that while low N_{IFN} and N_A are required for highly supercooled drizzle production (cf. Comstock et al., 2004; Hobbs & Rangno, 1998; Verlinde et al., 2013), large-scale ascent

may hasten and potentially increase the probability for the appearance of drizzle (e.g., Ikeda et al., 2007), but is not necessary for drizzle formation.

Figure 6 illustrates the LES output for the baseline model run at 9 h, compared with some of the observed parameters (see Appendix C for a description of the depicted parameter calculations). The model shows good agreement in the temperature (**Figure 6a**), q and the ice mass mixing ratio (q_i ; **Figure 6b**), RH_w and relative humidity with respect to ice (RH_i ; **Figure 6c**). The simulated cloud top as determined by the highest altitude of non-negligible q_c (**Figure 6b**) is also in good agreement with the observed CTH.

The contribution of turbulence, represented in **Figure 6f** by the total dissipation rate (ϵ_{tot}) that includes the contribution from numerics (see Chen et al., 2018), to the sharp temperature inversion just above cloud top (equivalent to the observed inversion; **Figure 6a**) was examined in other LES runs. These model runs in which turbulence did not develop (those with N_A of 10 cm^{-3} , see Figure S5 in the Supporting Information) or was prevented from developing (1-D runs; not shown) showed generally weaker inversions.

Examination of the drizzle and ice sedimentation fluxes at 1100 m, corresponding with the mean observed LCBH and the base of the main layer in the model, indicates that moisture removal from the cloud layer via hydrometeor sedimentation is dominated by drizzle (Figure S6). However, in the model runs where N_A and/or N_{IFN} are higher by a factor of 5, moisture sedimentation flux from the cloud layer is dominated by ice as a result of the limited q_r (see Figures S5 and S6). In these simulations, the lesser q_r relative to q_i also yields greater LDR below the LCBH (calculation description provided in Appendix C), which does not correspond as well with the observed profile as the model run results presented in **Figure 6e**.

Finally, the simulated cumulative optical depth (OD) profile shows some overestimation within the cloud boundaries (**Figure 6d**), although the total cloud OD in the model is in agreement with an estimation of its true value in the observed case (see Figure S6c in the Supporting information). In additional simulations initialized with a lower N_A , the cloud cumulative OD profile has a better agreement with the observations, because of the lower droplet number concentration. However, in these very small N_A simulations, the cloud-top radiative cooling is significantly reduced relative to the baseline simulation by lower LWP and cloud droplet number concentration (Figure S5; cf. Garrett & Zhao, 2006; Lubin & Vogelmann, 2006). Weaker cloud-top radiative cooling reduces turbulence production within the cloud layer, which is reflected in smaller ϵ_{tot} and a different cloud structure (e.g., q_c

distribution), all of which show poorer agreement with the observations (see Figure S5 in the Supporting Information).

Our results suggest that both N_{IFN} and N_A were likely limited to a range of rather small values. CCN number concentration measurements at the surface obtained at McMurdo Station during the drizzle event using a CCN counter (Uin, 2016) show mean and SD values that range from $8 \pm 3 \text{ cm}^{-3}$ at a supersaturation of 0.1% to $89 \pm 22 \text{ cm}^{-3}$ at 1%, respectively. However, direct conclusions regarding N_A aloft cannot be made in this drizzling event because the cloud layer is decoupled from the surface (**Figure 6a**), and thus aerosol concentration aloft is not linked to that at the surface (cf. Bromwich et al., 2012). With the absence of aerosol concentration measurements aloft, and based on the model comparison with the observations, we suggest that the activated N_{IFN} and N_A values during the drizzle event should have been on the order of 0.2 L^{-1} and 20 cm^{-3} , respectively. These values are within the typical range of McMurdo wintertime N_{IFN} (below $-20 \text{ }^\circ\text{C}$) and N_A from surface-based reports (Belosi et al., 2014; Liu et al., 2018). It is important to note that these number concentrations are rough estimates, because of numerous unexplored degrees of freedom in the model configuration as well as the many uncertainties associated with the comparison of the model results with the observations. Nonetheless, the use of plausible environmental conditions and parameter values yields model results generally consistent with available observations.

4. Summary and Conclusions

A persistent drizzle formation event at temperatures below -25°C was detected over McMurdo Station, Antarctica. From the Eulerian perspective of the observations, we showed that the supercooled drizzle may form at various LWP values and ice nucleation and growth rates. Large scale ascent of tens of centimeters per second may augment the drizzle formation, but short-term (with periodicities $< 60 \text{ min}$) gravity wave activity with the same or higher amplitudes do not show any apparent influence. This observational analysis shows that the August 16, 2016 drizzle formation event was possibly widespread along the Ross Ice Shelf coast.

Using large-eddy simulations, we demonstrated that persistent drizzle may form at these low temperatures regardless of the large-scale ascent magnitude, provided that the conditions are conducive to the formation of a sufficiently thick stratiform cloud layer, as long as the activated IFN and aerosol particle concentrations are sufficiently low. Under these relatively low aerosol concentration conditions, drizzle may comprise a substantial moisture sink from

a highly supercooled cloud layer, consistent with Mauritsen et al. (2011), and may be associated with the formation of multiple peaks in cloud water content profiles.

Persistent ascent that may hasten and/or increase the probability for drizzle formation is plausible in the orographically-forced atmosphere around the Antarctic coasts and above the inner continent (e.g., van den Broeke & van Lipzig, 2003; Parish & Bromwich, 1991).

Atmospheric temperatures below $-25\text{ }^{\circ}\text{C}$ and above the homogeneous freezing temperature regime are likely to occur near the coasts and inner continent around winter and summer months, respectively (e.g., Figure 7 in Comiso, 2000), while sufficiently low aerosol concentrations are possible year-round, mainly in the inner continent (e.g., Belosi et al., 2014; Bromwich et al., 2012; Shaw, 1988). Therefore, we posit that the necessary conditions for drizzle formation at these low temperatures (based on the observational and model analyses) can be met in the Antarctic continent. Moreover, these necessary conditions are likely to occur over the Southern Ocean as well, particularly during austral winter; Zhang et al. (2010) reported on the mixed-phase cloud temperatures below $-25\text{ }^{\circ}\text{C}$, while several studies documented low IFN and accumulation-mode aerosol concentrations (e.g., DeMott et al., 2018; Gras, 1995; Gras & Keywood, 2017; Schmale et al., 2019). While the literature lacks any reports of persistent drizzle formation at such low temperatures, we suggest that these persistent supercooled drizzle formation events may be rather common in Antarctica and the Southern Ocean. The spatial extent, climate model representation, and the potential impact of these events on the surface (e.g., radiative impact; see Mauritsen et al., 2011) are subjects for future research.

Appendix A: Estimate of Air Motion Bias in V_{dmax}

The effect of turbulent air motion is to broaden the fall spectrum of droplets within the radar volume (e.g., Doviak & Zrnić, 1993, Eq. 5.67; Gossard, 1988). V_{dmax} in the measured Doppler spectrum is where this broadened spectrum disappears into the spectrum noise floor (e.g., Rambukkange et al., 2011; Shupe et al., 2008). The offset of the spectrum edge from the actual fall speed of the drops depends not only on the turbulence but also on the radar operating characteristics that determine the radar volume, the radar minimum detectable signal, and the droplet population (e.g., Shupe et al., 2008)

We calculate the mean bias for the KAZR general (GE) mode data at 1200 m, 10-11 UTC (within the cloud layer; see **Figure 3**). The GE mode data are necessary for this calculation because of its sufficiently short radar pulse width. The ϵ_{tot} values used for the calculation

correspond with estimations following Section 4.1 in Chen et al. (2018), but with a smaller Gaussian filter width of 500 m, to account for the shorter time series.

The Doppler spectra are simulated using the forward model described in Appendix A in Chen et al. (2018), with the radar configurations taken from the KAZR deployed at McMurdo Station. For simplicity, we assume a mono-dispersed cloud droplet population, with LWC and droplet number concentration (N_d) estimates from the DHARMA simulations. For the dynamical broadening mechanisms, we only consider turbulent broadening and the beam width effect, with shear broadening neglected. The hydrometeor fall speed and vertical velocity are overlooked because we are only interested in the offset. The Doppler spectrum noise level (-94.6 dBm) for the 1200 m range gate is taken from the observed spectra, and background wind speed of 2 m/s is used based on the sounding measurements (**Figure 3h**).

We simulated 1000 realizations for each combination of the ε_{tot} , LWC, and number concentration of the cloud droplets to estimate the first Doppler spectral bin with reflectivity above the noise floor (**Figure A1**, top panels). An increase in the LWC and droplet diameters raises the offset by lifting the spectrum above the noise, whereas an increase in the dissipation rate broadens the spectrum, but lowers the SNR values. The mean offset values estimated varied from 0.05 m/s for low LWC (0.05 g/m^3) cases to over 0.3 m/s for LWC values of 0.2 g/m^3 , much larger than the 0.03 m/s reported by Shupe et al. (2008) for dynamical broadening. The calculated mean offset for the GE mode is approximately 0.1 m/s (for the DHARMA output parameters in the baseline run depicted in **Figure 6**), close to the observed 0.08 m/s GE mode mean V_{dmax} for the 1200 m range gate. We did not perform the full calculation for the higher sensitivity MD mode shown in **Figure 3** because of the complexity in defining the radar volume, but we did simulate the impact of the lower noise level in that mode (-98.4 dBm at 1200 m). The offset values based on the calculations with this noise level ranges (for the DHARMA output parameters) between ~0.2-0.3 m/s (**Figure A1**, bottom panels), which corresponds with the mean V_{dmax} value of 0.27 m/s at 1200 m in the MD measurements between 10-11 UTC.

Appendix B: DHARMA Model Setup

The following setup and physics options are used in DHARMA. A computational domain size of 3 km (x-axis) x 3 km (y-axis) x 10 km (z-axis) is used with a 50 m horizontal and 25 m vertical mesh between 500 and 2000 m, linearly stretched above and below. At each time step, the large-scale vertical motion profile is set to a uniform value between 100 and 3000 m

and interpolated linearly to zero at the surface below and at 4000 m aloft, where the uniform values are taken as the linearly interpolated mean BT ascent rates (black curve in **Figure 4b**). Dynamics are integrated with a maximum time step of 5 s, reduced when required to prevent the advective Courant-Friedrichs-Lewy number from exceeding 0.8. Microphysics is treated following the Morrison et al. (2009) two-moment scheme with cloud water, rain, and cloud ice (no snow or dense ice class), with a prognostic saturation excess following Morrison and Grabowski (2008). The relative dispersion of the cloud drop size distribution is fixed at 0.3 per Geoffroy et al. (2010). Diurnal radiative transfer is computed column-wise every 60 s in 44 wavelength bins using a two-stream model (Toon et al., 1989) with water vapor continuum absorption (Clough et al., 1989). The temperature and water vapor of the model overlying the LES model domain, which control the downwelling longwave flux at the domain top are adjusted to match fluxes derived from observations. Liquid particle scattering and absorption coefficients are computed from the Lorenz-Mie theory (Toon & Ackerman, 1981). Ice particle scattering and absorption coefficients account for non-sphericity using the approach described by van Diedenhoven et al. (2012). Surface skin temperature is fixed at 250 K, and surface turbulent fluxes are fixed to zero.

For simplicity, we specify a monomodal lognormal aerosol particle size distribution with a geometric mean radius of 0.1 μm and geometric SD of 2, approximated from surface measurements during AWARE (Liu et al., 2018). Aerosol activation follows Abdul-Razzak and Ghan (2000), in which the supersaturation is taken as the local minimum during a time step using the semi-analytic treatment of Morrison and Grabowski (2008), which neglects activation (cf. Vogelmann et al., 2015). Ice initiation is limited to idealized immersion-mode freezing using the approach described by Ovchinnikov et al. (2014), wherein cloud ice crystals are introduced wherever a grid cell is liquid-saturated and the ice number concentration is below the specified value, referred to as N_{IFN} . This approach leaves two variable parameters: the number concentrations of the aerosol particle size distribution (N_A) and the threshold in-cloud ice number concentration below which additional ice will be initiated (N_{IFN}).

Appendix C: Retrieval of the Observational Parameters and Modeled LDR

The 10:41 UTC radiosonde measurements are used as the temperature and RH_w profiles. Measurement uncertainty values are taken from Holdridge et al. (2011). q and RH_i are calculated using the formulation in Murphy and Koop (2005).

The q_i is calculated by dividing the IWC estimates by the dry air density (calculated based on the sounding measurements). The IWC is estimated from the radar reflectivity and the temperature profile following Hogan et al. (2006). The q_i uncertainty is taken from the potential bias in the methodology as shown in Heymsfield et al. (2008, their Figures 9-15).

The dissipation rate ε_{tot} is calculated between 1020 m and 1470 m in 90 m intervals using the KAZR's GE mode data, following the methodology of Chen et al. (2018). ε_{tot} is not calculated for higher altitudes, because continuous KAZR records are necessary, and the GE mode data are rather patchy closer to cloud top. The ε_{tot} calculations are performed over the longest continuous KAZR record in each altitude between 10-11 UTC, using a Gaussian filter width of 0.5 km. The ε_{tot} uncertainty is estimated to be 60%, based on the ε_{tot} sensitivity to the width of the Gaussian filter (see Chen et al., 2018).

The observed drizzle base altitude is determined 7.5 m (one HSRL range gate) above the highest HSRL range gate where $\beta_p < 10^{-5.3} \text{ m}^{-1} \text{ sr}^{-1}$ and the LDR ≥ 0.2 . These β_p and LDR threshold values were determined after careful examination of the HSRL parameters together with the KAZR spectra throughout the drizzle event. The values of the observed drizzle base altitude, LCBH, CTH, particulate optical depth (OD_{tot} ; taken from the HSRL measurements), and LDR, and their uncertainties (all illustrated in **Figure 6**) are taken as the 10-11 UTC hourly mean and SD, respectively.

The model LDR is calculated at each altitude level by weighting the typical LDR values for liquid and ice by the horizontal domain's mean projected area per unit volume of liquid (droplets + rain) and ice hydrometeors, respectively (a hydrometeor population's projected area is excluded when its mass mixing ratio is smaller than 10^{-3} g/km). The typical LDR values for liquid (0.025) and ice (0.32) are chosen as the LDR value of these hydrometeor population peaks in a 2-D β_p versus LDR histogram (e.g., Figure 2 in Silber et al., 2018a), based on the 10-11 UTC hourly data.

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<https://cds.climate.copernicus.eu>) and the ECMWF database (<https://apps.ecmwf.int/datasets/data/interim-full-daily/>), respectively. CERES and CALIOP observational data can be downloaded from the NASA Langley Research Center website (https://ceres.larc.nasa.gov/order_data.php) and the NASA Atmospheric Science Data Center website (https://eosweb.larc.nasa.gov/project/calipso/calipso_table), respectively. AMPS data can be requested from the Ohio State University Polar Meteorology Group (<http://polarmet.osu.edu/AMPS/>). Wavelet software was provided by C. Torrence and G. Compo, and is available at: <http://paos.colorado.edu/research/wavelets/>. The authors wish to thank Paul Demott for the fruitful discussion. I.S. and J.V. are supported by the National Science Foundation grant PLR-1443495 and by the DOE grant DE-SC0017981. I.S. is also supported by DOE grant DE-SC0018046. A.F. and A.A. are supported by the NASA Radiation Science and Modeling, Analysis and Prediction programs. D.H.B. and S.H.W. are supported by National Science Foundation grant PLR-1443443 and by DOE grant DE-SC0017981. Contribution XXXX of Byrd Polar & Climate Research Center.

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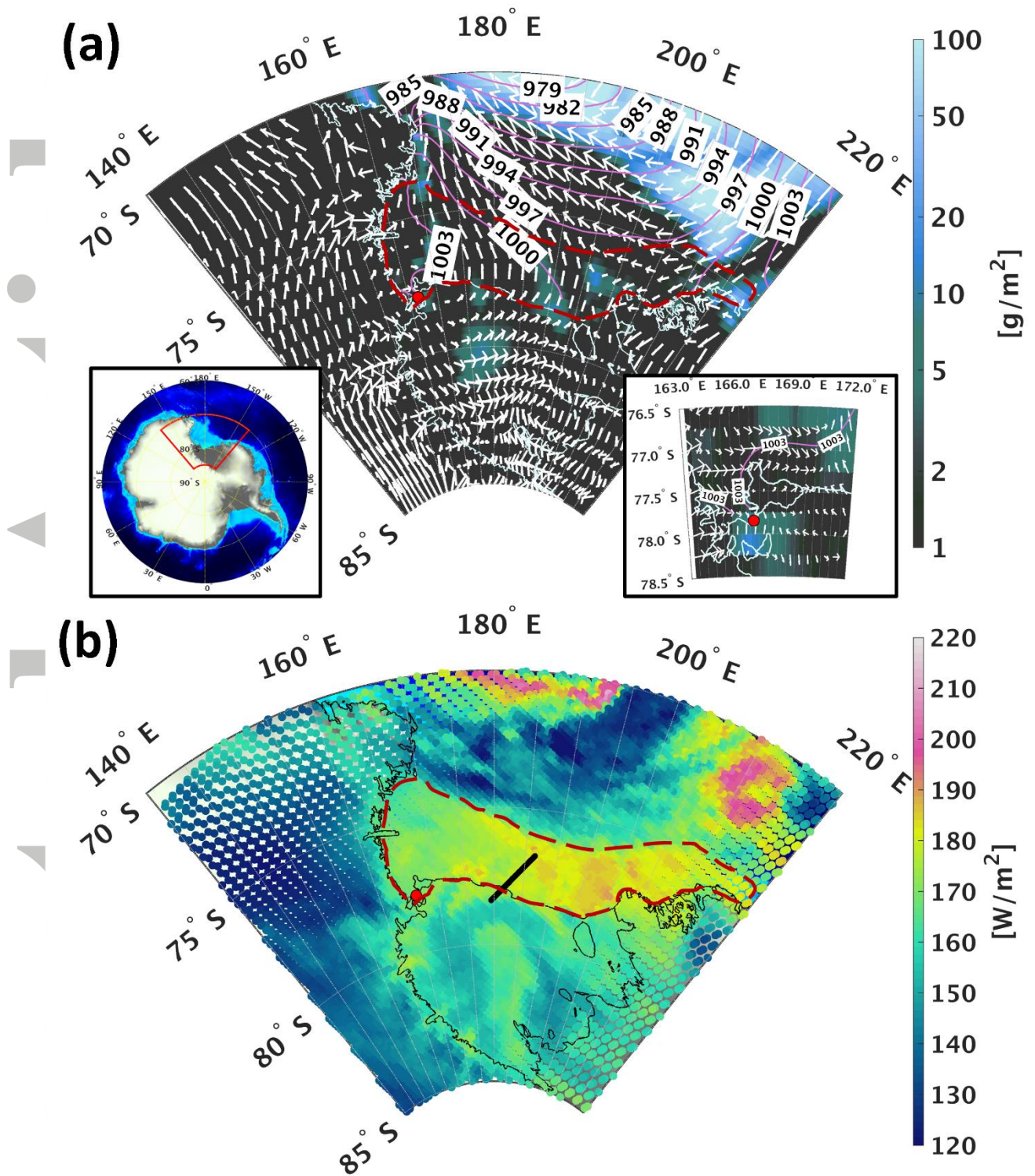


Figure 1: (a) Sea level pressure (SLP; in hPa; land areas are masked), integrated liquid water (cloud + rain; log-scaled color bar), and surface winds illustrated in contours, color-scale, and quivers, respectively, resolved by ERA5 for August 16, 2016 at 10:00-11:00 UTC. The red circle marker designates the location of McMurdo Station (located on Ross Island). The inset panels show a topographic map of Antarctica (the red box

designates the region depicted in the main panels) and a focus on the flow field surrounding Ross Island. (b) Top of atmosphere upwelling longwave radiation footprints measured by CERES Aqua on August 16, 2016, between 10:44 and 10:50 UTC. The dashed red shape highlights the region of enhanced upwelling longwave radiation discussed in the text. The thick black curve designates part of the CALIPSO satellite trajectory (10:48-10:50 UTC), in which low-level mixed-phase clouds were detected (see Figure S2 in the Supporting Information). The 1 arc-minute topographic data were developed by the National Geophysical Data Center (NGDC, Amante, 2009), and is freely available at <https://www.ngdc.noaa.gov/mgg/global/global.html>.

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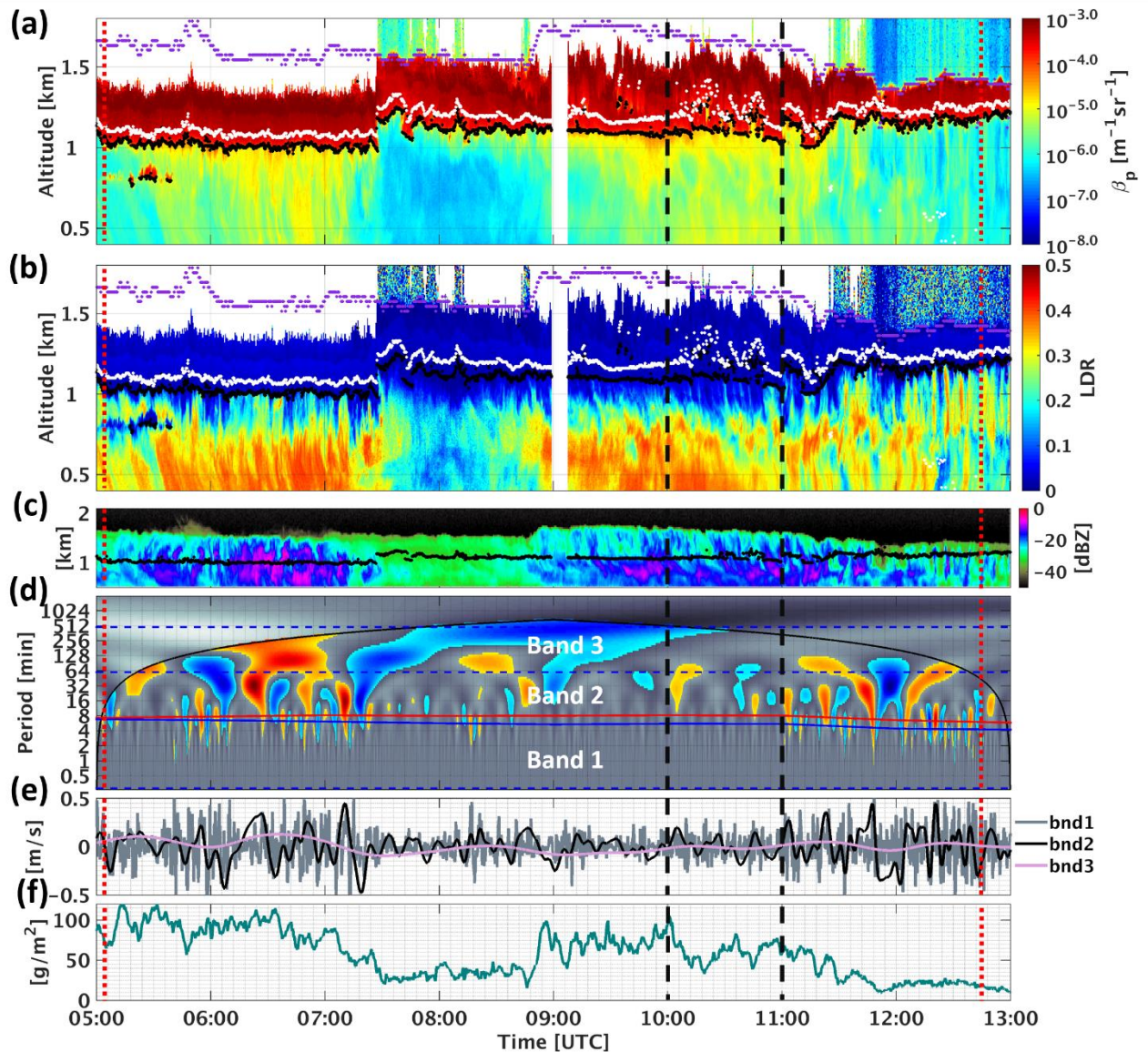


Figure 2: (a) High Spectral Resolution Lidar (HSRL) particulate backscatter cross-section (β_p ; color bar in log-scale). The white and black dots show the ceilometer cloud and HSRL liquid cloud base height (LCBH) data product outputs, respectively. Purple dots indicate the cloud-top height (CTH) determined using the Ka-band ARM Zenith Radar (KAZR) moderate sensitivity (MD) mode measurements (see text for details). Note that the apparent jump in the observed liquid-bearing cloud layer at 07:27 UTC is the result of a (physical) rapid (~ 1 min) ~ 200 m upward shift in the lowest LCBH. (b) HSRL linear depolarization ratio (LDR). (c) KAZR integrated reflectivity factor. The “main” HSRL LCBHs used for the wavelet analysis are designated by the black dots. (d) wavelet analysis of the most upward Doppler velocity in the KAZR Doppler velocity spectrum with signal above the noise floor (V_{dmax}) 100 m above the main LCBH, which provides an estimate of the vertical air motion. The colored regions denote statistically significant periodicities at the 99% level. The solid black curve represents

the wavelet cone of influence. The solid red and blue curves indicate the Brunt-Vaisala Frequency (BVF) and the Eulerian apparent BVF for a horizontal wavelength of 4 km, respectively, calculated for McMurdo Station based on the atmospheric parameters 0-200 m above the main LCBH using the 1.1 km AMPS model output (winds are taken at the LCBH + 100 m). The solid blue curve (apparent BVF) and the 3 horizontal dashed blue lines (located at 17 s, 60 min, and 8 h) partition the spectrogram into 3 main bands, which can represent (from the smallest to largest time scale) the cloud-scale circulation (and inhomogeneity), gravity wave scale oscillations, and synoptic-scale air motion. (e) Temporal variability of the three bands (see legend for details). (f) liquid water path derived from microwave radiometer measurements. The vertical dotted red lines define the beginning and ending of the drizzle event. The vertical dashed black lines indicate the edges of the period represented in Figure 1 and in Figure 3.

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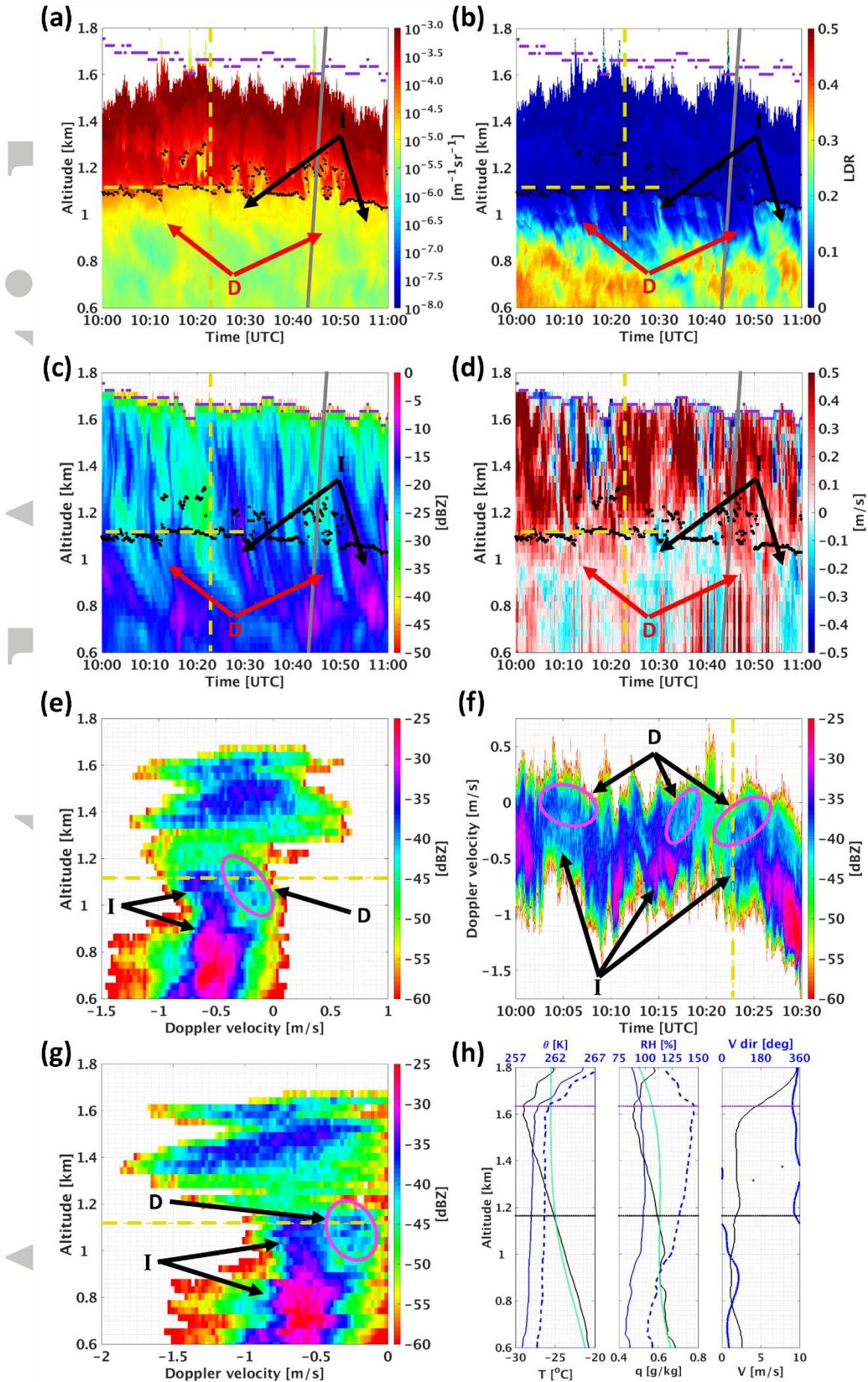


Figure 3: Measurements gathered during the drizzle event on August 16, 2016, between 10:00-11:00 UTC: (a) HSRL β_p (log-scaled colorbar), (b) HSRL LDR, (c) KAZR integrated reflectivity factor, (d) KAZR V_{dmax} , (e) vertical profile of the KAZR Doppler spectra at 10:22:47 UTC, (f) temporal evolution of the KAZR Doppler spectra at 1118 m altitude between 10:00-10:30 UTC, (g) same as panel e, but with the spectra shifted by V_{dmax} to 0 m/s, (h) balloon (released at 10:41 UTC) sounding measurement profiles, where the solid black curves in the left, middle, and right panels illustrate the temperature, specific humidity (q), and horizontal wind speed, and the solid (dashed) blue curve shows the potential temperature θ (equivalent potential temperature θ_e), relative humidity with respect to water (ice), and the wind direction, respectively. The dotted green curves in the left and middle sounding panels depict the temperature and q profiles (respectively) used for the 1-D “backward” model run (see Section 3). All the KAZR parameters shown in this figure are from the MD mode measurements. The gray lines illustrated in panels a-d denote the balloon time-altitude location (horizontal coordinates are different from the HSRL and KAZR field of view). The black and purple dots in panels a-d and h mark the HSRL liquid cloud base height and KAZR cloud-top height, respectively. The vertical (horizontal) dashed yellow lines in panels a-g designate the time (altitude and time period) of the spectra presented in panels e and g (f). The letter ‘D’ and its corresponding arrows and ellipses point at some of the drizzle time-height locations in the HSRL and KAZR profiles (panels a-d), and indicate the drizzle spectra signatures in panels e-g (ice crystal locations and spectra signatures are indicated by the letter ‘I’).

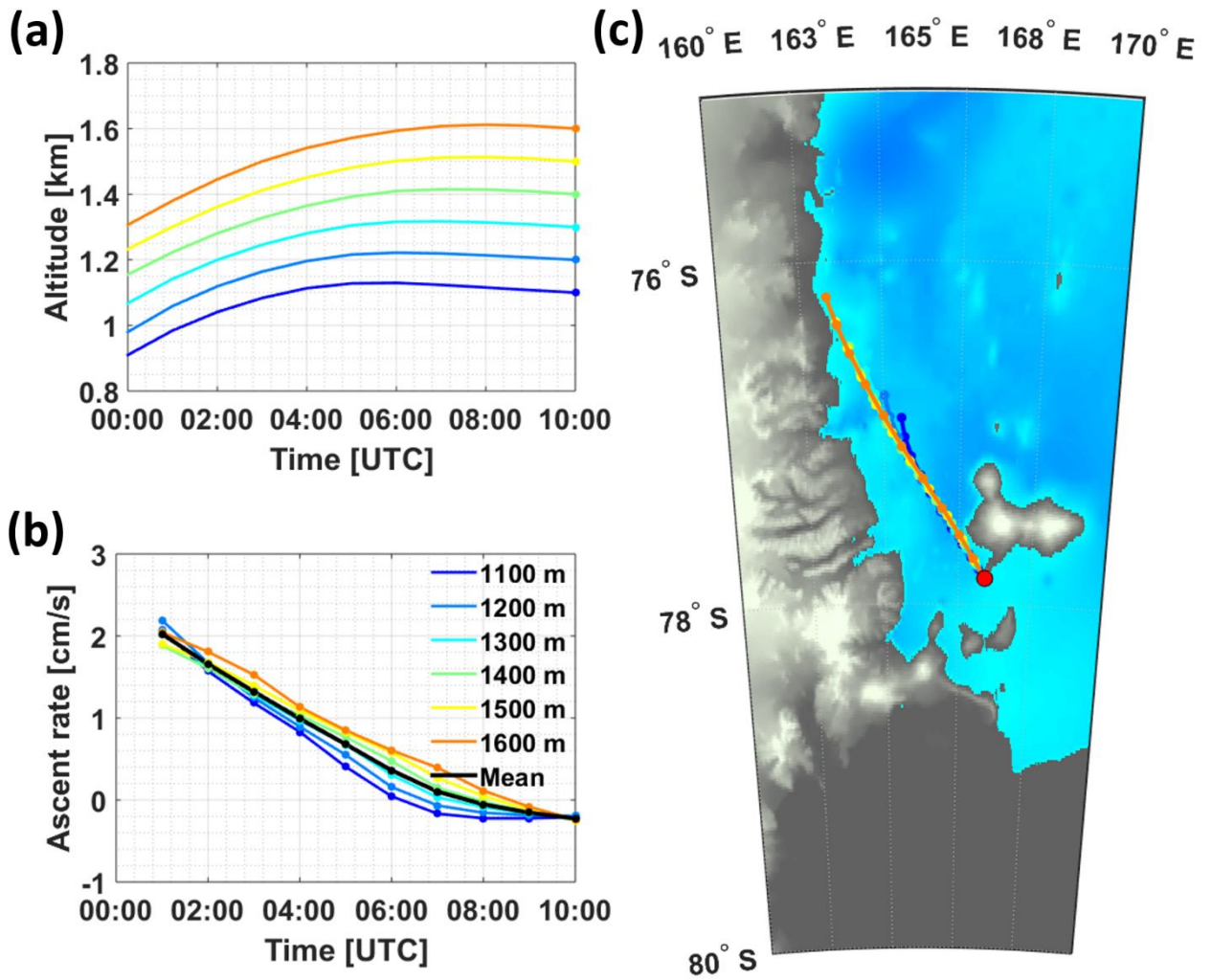


Figure 4: 10 h Back trajectories and 9 h mean ascent rates of air parcels at cloud layer heights (following the observations shown in Figure 3; see legend in panel b), based on the HYSPLIT model initialized with ERA-Interim reanalysis data (see text for details). The air parcel altitude, hourly ascent rates, and coordinates are shown in panels a, b, and c, respectively. The red circle marker in panel c designates the location of McMurdo Station.

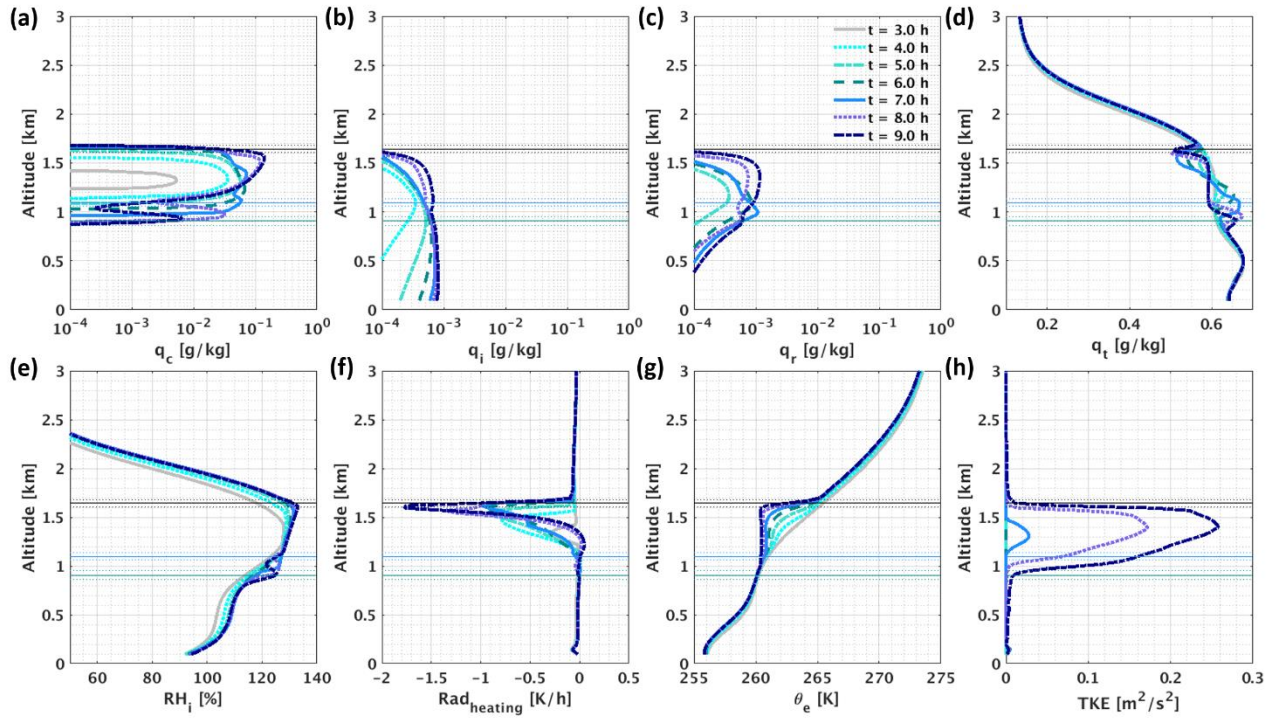


Figure 5: Domain-average profiles of DHARMA model output at different times (see legend) for ice freezing nuclei and accumulation-mode aerosol particle number concentrations of 0.2 L^{-1} and 20 cm^{-3} , respectively. (a-d) Mass mixing ratios of cloud water (q_c), cloud ice (q_i), rain (q_r), and total water (q_t), respectively, (e) relative humidity with respect to ice (RH_i), (f) radiative heating rate, (g) θ_e , and (h) turbulent kinetic energy (TKE). The solid (dotted) horizontal black, blue, and green lines represent the observed 10:00-11:00 UTC hourly mean (mean \pm standard deviation) CTH, LCBH, and drizzle base height, respectively.

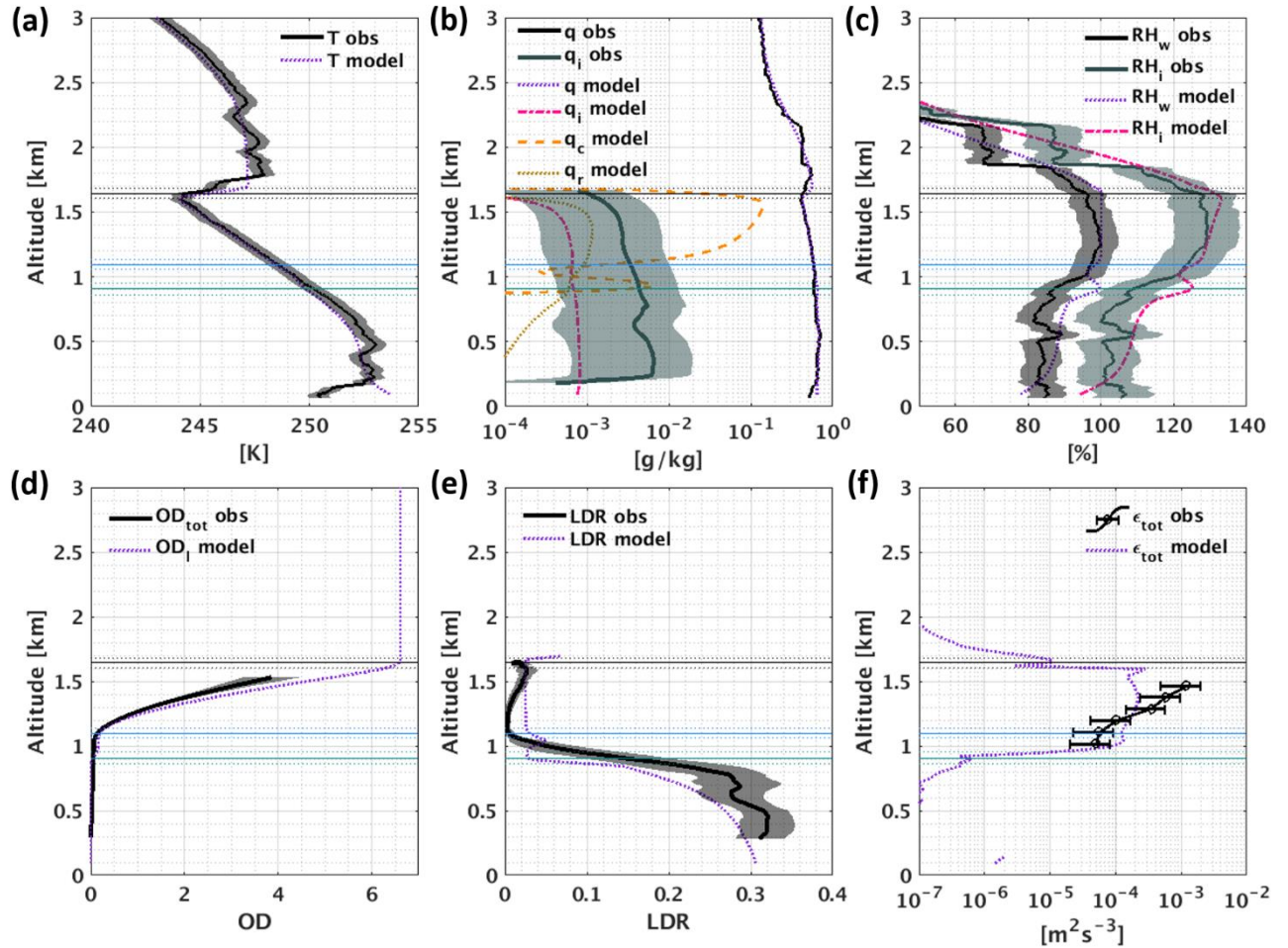


Figure 6: Observed and DHARMA model output domain-average profiles (model initialized using ice freezing nuclei and accumulation-mode aerosol particle number concentrations of 0.2 L^{-1} and 20 cm^{-3}). (a) Temperature, (b) q , q_i , q_c , and q_r , (c) relative humidity with respect to water (RH_w) and RH_i , (d) cumulative particulate and liquid water optical depths (OD_{tot} and OD_i , respectively), (e) LDR, and (f) total dissipation rate (ϵ_{tot}). The shaded areas (panels a-e) and error bars (panel f) designate the observation uncertainty. The solid and dotted horizontal lines are the same in Figure 5. The faint HSRL signal near the CTH allows the derivation of the LDR (with high uncertainty) but not the calculation of the OD_{tot} . For a description of the observational parameter retrievals and the modeled LDR calculation see Appendix C.

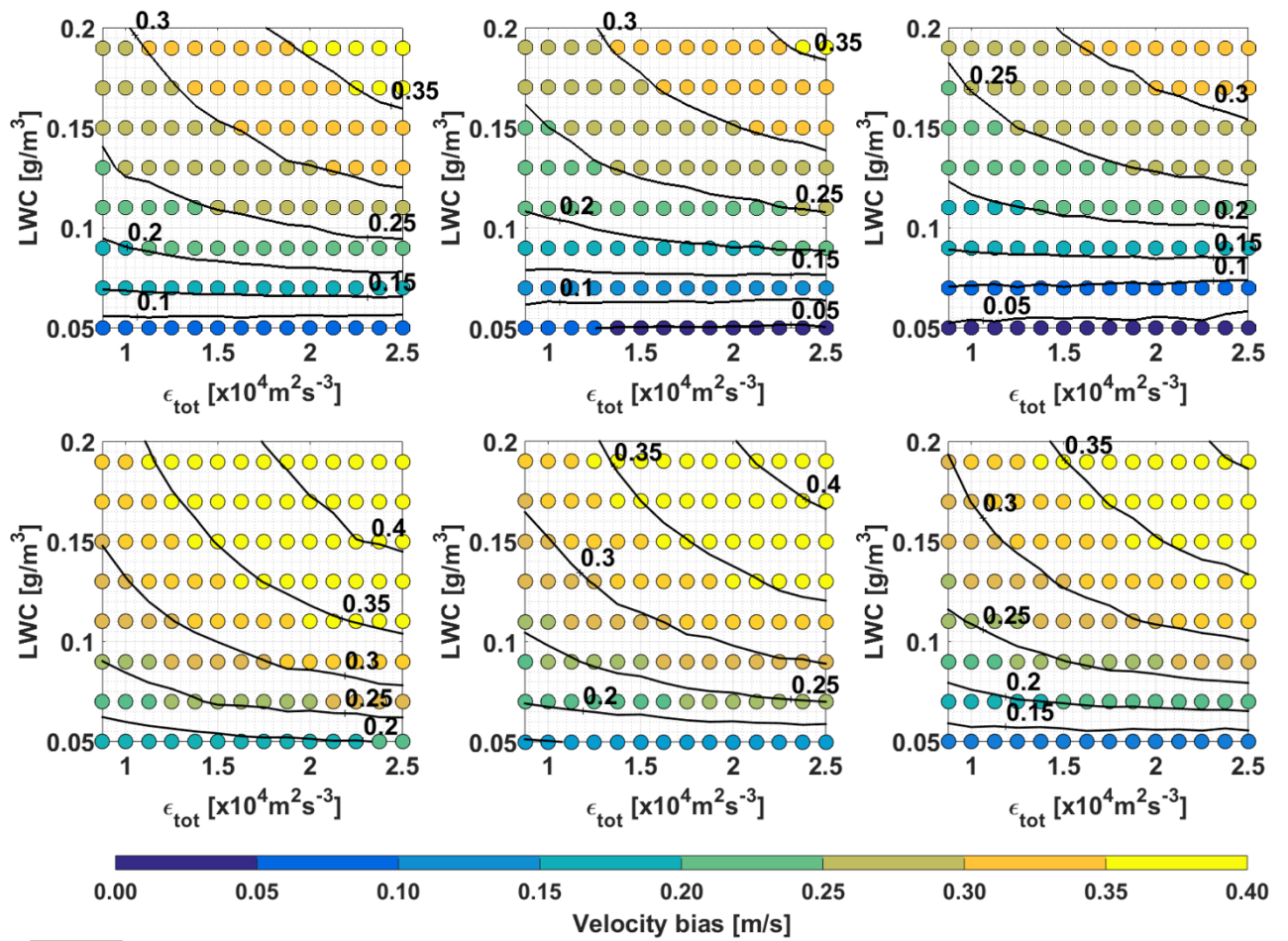


Figure A1: Mean bias (in m/s) of the true air velocity represented in the KAZR general (GE) mode V_{dmax} for this case (top). The panels represent (from left to right) increasing droplet number concentration (15, 20, and 25 cm⁻³, respectively). Each circle marker represents the mean bias for one set of 1000 simulations and the contours are isopleths of offset to indicate the dependencies. The bottom panels show the results after using the noise level in KAZR MD mode.