1. Introduction

When a molecular gas is condensed onto a substrate, a solid film may be formed which spontaneously exhibits a static electric field,\(^1\,^\text{--}\,^8\) the strength of which may exceed \(10^8\) V m\(^{-1}\). Such spontaneously electrical solid films, so-called ‘spontelectrics’, represent a new structural and electrical phase of the solid state. The present work is devoted to a demonstration that RAIRS may be used to establish new spontelectric materials. This is illustrated here for solid carbon monoxide between 20 and 26 K.

The proof of principle of this method was provided recently through a RAIRS study of \(\text{N}_2\text{O}\) films.\(^8\) Such films are known to be spontelectric,\(^1\) through direct measurement of surface potentials, using an electron beam technique. Results in ref. 8 show that the spontelectric field may be detected by observing the vibrational Stark effect in the material, using RAIRS. The known strength and temperature dependence of the field in \(\text{N}_2\text{O}\) were used to reproduce the corresponding temperature dependence of the longitudinal–transverse optical (LO–TO) splitting. Here we invert the argument and use temperature dependence of the LO–TO splitting to obtain the spontelectric field in CO. The crux of the technique is that a measurable temperature dependence of the LO–TO splitting is diagnostic of the spontlectric nature of the film of material concerned.

On this basis, we use RAIRS to estimate the spontlectric parameters of solid CO, for which, in contrast to \(\text{N}_2\text{O}\), no direct measurements of surface potentials have been made. In this study, we accordingly find that the LO–TO splitting in RAIR spectra of CO depends on the temperature of deposition, to the tune of \(\sim 0.045\) cm\(^{-1}\) K\(^{-1}\). Measurements of the LO–TO splitting then allow the determination of the values of the spontlectric field as a function of film deposition temperature.

The reasons for pursuing this work are: (i) RAIRS is a technique available to many laboratories and RAIRS could therefore be used to investigate the possible spontlectric nature of many materials. (ii) There are potentially important astrophysical implications if CO is spontlectric.

Spontelectric materials discovered to date all possess a permanent dipole moment and range over simple hydrocarbons, halocarbons, alcohols, organic formates, benzene derivatives and such simple inorganics as nitrous oxide. Numerous data, outlined in ref. 1 and 7, and corresponding analyses, point to dipole orientation as the origin of the spontlectric polarization. This in turn gives rise to the electric fields observed. The salient properties of spontelectrics\(^1,^6,^7\) are that (i) the spontlectric field is lower for higher deposition temperatures [but see \(5\)], (ii) the nature of the substrate, on which materials are condensed, has no bearing on the strength of the bulk spontlectric field, (iii) the value of the spontelectric field depends on both the nature of the material which is deposited and on the temperature at which the film is deposited, (iv) there exists a critical annealing temperature, termed the Curie point, by analogy with ferromagnetism, at which films depolarize and the spontlectric field disappears.

The presence of the spontlectric field results in a vibrational Stark effect in the solid, causing a shift in characteristic
vibrational frequencies.\textsuperscript{8–12} Since the strength of the spontaneous electric field depends strongly on the temperature of deposition of the film of material, there is a corresponding temperature dependence of the vibrational frequencies measured using RAIRS. Based upon a model for the spontaneous effect,\textsuperscript{1} we show in Section 3 that the apparent LO–TO splitting in solid CO has a significant contribution from the Stark effect arising from the spontaneous electric field. We find, for example, that at a deposition temperature of 20 K, the Stark effect contributes \( \sim 36\% \) of the total measured splitting of 4.03 cm\(^{-1}\). The observed temperature dependence of LO–TO splitting may be attributed wholly to the dependence of the spontaneous field on the film deposition temperature.\textsuperscript{8}

For simplicity, we refer throughout to the observed splitting in RAIR spectra of CO films as LO–TO splitting. At the same time, we recognize that the absolute value of the splitting arises through a combination of the intrinsically different vibrational frequencies associated with LO and TO modes and, at the level of approximation adopted here, an independent contribution due to the vibrational Stark effect.

In the current work, spontaneous films are interrogated using RAIRS with a grazing infrared beam, such that the incident electric field of the beam has components both parallel and perpendicular to the film normal. Relative to the incident beam wavelength, the film can be considered infinite in the plane of the film and only transverse optical (TO) phonons can be excited in this plane. If however the thickness of the film is comparable to the wavelength of the incident beam, the boundary conditions allow for the excitation of longitudinal optical (LO) phonons along the normal axis. This is known as the Berreman effect,\textsuperscript{13} and has been studied extensively in non-ionic films.\textsuperscript{14–16} Longitudinal phonons resonate at higher frequencies, because of the induced field associated with longitudinal waves passing through a medium composed of dipolar species. Thus, LO–TO splitting occurs for vibrational modes, when an incident beam interrogates a thin film at a suitably oblique angle.

Henceforth \( \nu_L \) and \( \nu_T \) represent the frequencies for LO and TO phonons respectively and \( \Delta \nu = \nu_L - \nu_T \) represents the value of the splitting. The force fields giving rise to \( \nu_L \) and \( \nu_T \) are modified by the vibrational Stark effect, through the presence of the spontaneous field. We show in Section 3 how we may relate the resulting modification of LO–TO splitting, and its temperature dependence, to the presence of a static spontaneous electric field, oriented along the surface normal of the film. Our analysis demonstrates how measurements of \( \nu_L \) and \( \nu_T \) may then lead to a full characterization of the spontaneous field in CO.

2. Experimental method and results

2.1 The experimental method

RAIRS experiments were performed in an ultrahigh vacuum system, described in detail elsewhere.\textsuperscript{7,14} The substrate, an oxygen-free high conductivity copper block coated with a 300 nm amorphous silica layer,\textsuperscript{19} is mounted on the end of a closed-cycle helium cryostat, reaching a base temperature of 18 K, measured with a KP-type thermocouple connected to an IJ-6 temperature controller (IJ Instruments). The central chamber is equipped with a line-of-sight quadrupole mass spectrometer (QMS, Hiden Analytical) and a Fourier-transform infrared spectrometer (Varian 670-IR) used in reflection–absorption mode, at a grazing incidence of 75\(^{\circ}\) with respect to the normal to the substrate. After reflection from the sample, the infrared beam is focused into a liquid nitrogen cooled HgCdTe detector. The RAIR spectra presented here result from the co-addition of 512 spectra recorded at 0.1 cm\(^{-1}\) resolution. This high resolution was necessary to measure the small frequency shifts observed in RAIR spectra of solid CO films, recorded at different deposition temperatures.

Films are deposited by background dosing of CO gas (BOC, purity 99.9\%) onto the substrate at a rate of 0.05 ML s\(^{-1}\). Thicknesses of CO films in monolayers (ML) were determined (\( \pm 20\% \)) through temperature-programmed desorption experiments, performed by applying a heating ramp of 0.3 K s\(^{-1}\) from the deposition temperature, with the desorbed species detected using the QMS.

The choice of the substrate was determined by the metal surface selection rule, which dictates that TO modes are silent on a metal surface. The presence of the silica layer coating, on the copper, relaxes this selection rule and allows the observation of both LO and TO modes in solid CO on silica, while retaining the enhanced sensitivity associated with RAIR spectroscopy. For a more detailed discussion in the case of N\(_2\)O films, see ref. 8.

2.2 Results

Fig. 1 presents the RAIR spectra of the \( \nu_{CO} \) band of 5 ML CO films, deposited on 300 nm silica between 20 and 26 K. Increasing the deposition temperature red-shifts the LO mode, whilst the TO mode is blue-shifted. Annealing the films from 18 K to 26 K (not shown) does not produce any detectable shift of the bands. During the annealing experiments, desorption was
negligible at and below 26 K, as confirmed by TPD and by the almost constant intensity of the RAIR signal of films annealed from 20 to 26 K.

LO–TO splitting can be observed at lower temperature than in the data shown here, that is, below 20 K. Evidence for volcano desorption from our substrate, during TPD of N₂O trapped in CO, is consistent with CO undergoing a phase change at around 20 K. Here RAIRS data are only presented for crystalline CO films, that is, for deposition temperatures ≥ 20 K, and the present study is therefore limited to a discussion of such crystalline films. The upper temperature boundary is set by the observation that films of CO form a steady state below 26 K, where the rate of desorption is low. Above 26 K the rate of desorption increases such that stable films cannot be grown.²⁰

In the deposition experiments, the νCO LO and TO modes were fitted with Gaussian functions using the Igor Pro software. Fig. 2 presents the RAIR spectrum of a 5 ML CO film deposited at 20 K on 300 nm SiO₂ (open symbols), and the two Gaussian functions giving the best agreement with the experimental LO and TO modes (full lines). The inset of Fig. 2 shows the residuals, that is, the difference between the experimental spectrum and the fitting function, obtained in the LO mode region, with the best fit (full line) and after variation of the central position of the Gaussian by +0.01 (dotted line) and −0.01 cm⁻¹ (dashed line). One can see that the residual is almost zero in the central region, and increases when going away from the maximum because of the inhomogeneous broadening of the bands. The flattest residual is obtained for a central position of 2142.52 cm⁻¹, that was hence considered the best fit for determination of the LO frequency at 20 K (see Table 2). A similar procedure was used for all determinations of LO and TO frequencies presented here. The uncertainties quoted above correspond to the maximum variation that can be applied to the central value of the fitted peak whilst maintaining the best match between the experimental spectrum and the fitted curve.

This method gives access to the uncertainty in the band position, yielding ±0.01 cm⁻¹ for LO modes (see above and inset to Fig. 2) and ±0.02 cm⁻¹ for the broader TO modes. The band positions resulting from the fits are displayed in Fig. 3 with stars and circles for the νCO LO and TO modes, respectively, for 5 ML CO films, deposited on 300 nm silica, as a function of deposition temperature.

We may also use the variation of the inhomogeneous broadening of LO and TO bands, with deposition temperature, as a qualitative indication of the degree of dipole orientation in the film. Fig. 4 shows a measurement of the inhomogeneous broadening of the νCO band for different deposition temperatures. The degree of broadening is estimated here by measuring the intensity of the RAIR spectrum at 2141 cm⁻¹, essentially the average frequency of the LO and TO modes, and normalizing by the integrated area of the band. This allows comparison
between different experiments and yields the ordinate in Fig. 4. Insofar as inhomogeneous broadening is a measure of the range of environments in which any component species finds itself, the greater the inhomogeneous broadening the less the dipole orientation. Results in Fig. 4 are therefore consistent with the expected behavior of a spontelectric material, and may be understood to show that the drop in dipole orientation increases with deposition temperature, as already observed in N$_2$O and numerous other films.1,8

3. A model for the spontelectric Stark effect

We first present qualitative evidence that CO films are indeed spontelectric. The LO–TO splitting changes from 4.02 cm$^{-1}$ in 20 K films to 3.75 cm$^{-1}$ at 26 K, noting that we have mentioned in the introduction that such variation is diagnostic of the presence of a temperature dependent spontelectric field in the film.

In ref. 8, we considered in detail for N$_2$O whether thermal expansion$^{14}$ could lead to the observed temperature variation of LO–TO splitting, $\Delta\nu$. Here, the measured average value of $d\Delta\nu/dT \sim -0.045$ cm$^{-1}$ K$^{-1}$ (see Table 2). Assuming that figures of density for solid CO are similar to those for N$_2$O, this would require a volume expansion coefficient for CO of between 0.002 and 0.003 K$^{-1}$, using the same analysis as that presented in Section 3.1.1 of ref. 8. However, low temperature volume expansion coefficients for such solid molecular materials typically lie between $10^{-5}$ and $10^{-6}$ K$^{-1}$. Therefore thermal expansion effects have a negligibly small influence on $\Delta\nu$. We conclude that the CO data presented here are prima facie evidence for the spontelectric nature of CO films.

We now seek to establish the spontelectric parameters of solid CO as a function of deposition temperature from the experimental data presented in Section 2. The most significant quantities, emerging from the subsequent analysis, are the values of the spontelectric field and the accompanying degrees of dipole orientation and of bulk polarization, as a function of the temperature of deposition.

In order to proceed, we first present our data in tabular form in Table 2 in Section 3.1. We use these data to evaluate the ratio of the Stark splitting contribution, $\Delta\nu_{\text{St}}$, to the full LO–TO splitting, $\Delta\nu$. In Section 3.2 we briefly review the necessary theoretical background for spontelectrics, identifying the symmetric field and asymmetric field parameters, $\langle E_{\text{sym}} \rangle$ and $\langle E_{\text{asym}} \rangle$ respectively, and defining the degree of dipole orientation. In Section 3.3, we develop explicit expressions relating $\Delta\nu_{\text{St}}$ to the spontelectric parameters and in Section 3.4 we extract values of the spontelectric field as a function of deposition temperature. For ease of reference, Table 1 shows the definition of the various symbols used here.

### Table 1
<table>
<thead>
<tr>
<th>Symbol</th>
<th>Description</th>
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<tbody>
<tr>
<td>$\nu_L$</td>
<td>Longitudinal optical (LO) frequency</td>
</tr>
<tr>
<td>$\nu_T$</td>
<td>Transverse optical (TO) frequency</td>
</tr>
<tr>
<td>$\Delta\nu$</td>
<td>Measured LO–TO splitting</td>
</tr>
<tr>
<td>$\Delta\nu_{\text{St}}$</td>
<td>Splitting due to the spontelectric Stark field</td>
</tr>
<tr>
<td>$\Delta\nu_{\text{B}}$</td>
<td>Intrinsic splitting</td>
</tr>
<tr>
<td>$\xi$</td>
<td>$\Delta\nu_{\text{St}}/\Delta\nu$</td>
</tr>
<tr>
<td>$\mu$</td>
<td>Dipole moment of CO in the solid state</td>
</tr>
<tr>
<td>$\mu_{\text{sym}}$, $\mu_{\text{asym}}$</td>
<td>Symmetric field parameter, Asymmetric field parameter</td>
</tr>
<tr>
<td>$\langle E_{\text{sym}} \rangle$, $\langle E_{\text{asym}} \rangle$</td>
<td>$\langle \mu \rangle \mu / \langle \mu \rangle O \Omega$</td>
</tr>
</tbody>
</table>

### Table 2
<table>
<thead>
<tr>
<th>$T/K$</th>
<th>$\nu_L/\text{cm}^{-1}$ $\pm 0.02$</th>
<th>$\nu_T/\text{cm}^{-1}$ $\pm 0.01$</th>
<th>$\Delta\nu/\text{cm}^{-1}$</th>
<th>$\Delta\nu_{\text{St}}/\text{cm}^{-1}$</th>
<th>$\xi$</th>
</tr>
</thead>
<tbody>
<tr>
<td>20</td>
<td>2138.50</td>
<td>2142.52</td>
<td>4.02</td>
<td>1.44</td>
<td>0.357</td>
</tr>
<tr>
<td>21</td>
<td>2138.55</td>
<td>2142.51</td>
<td>3.96</td>
<td>1.37</td>
<td>0.345</td>
</tr>
<tr>
<td>22</td>
<td>2138.58</td>
<td>2142.48</td>
<td>3.90</td>
<td>1.31</td>
<td>0.335</td>
</tr>
<tr>
<td>24</td>
<td>2138.62</td>
<td>2142.41</td>
<td>3.79</td>
<td>1.20</td>
<td>0.316</td>
</tr>
<tr>
<td>26</td>
<td>2138.63</td>
<td>2142.38</td>
<td>3.75</td>
<td>1.16</td>
<td>0.308</td>
</tr>
</tbody>
</table>

data, for LO–TO frequencies in solid CO, as a function of deposition temperature, $T$. These data are shown in Fig. 3. Column 4 of Table 1 shows the LO–TO splitting, $\Delta\nu_{\text{B}}$, that is, column 3–column 2.

$\Delta\nu$ may be represented by the sum of two terms, one of which, the intrinsic splitting, $\Delta\nu_{\text{B}}$, is independent of temperature of deposition, whereas the other, $\Delta\nu_{\text{St}}$, the spontelectric term, is a function of deposition temperature. Thus at any temperature $\Delta\nu = \Delta\nu_{\text{B}} + \Delta\nu_{\text{St}}$. We can use the data in Table 2 to determine an experimentally based value of $\Delta\nu_{\text{B}}$. This gives the desired values of $\Delta\nu_{\text{St}}$, as a function of temperature.

In order to estimate the relative contributions of $\Delta\nu_{\text{St}}$ and $\Delta\nu_{\text{B}}$ as a function of temperature, we invoke the concept that the spontelectric Stark field must tend to zero as the temperature is indefinitely raised, since at some temperature the system encounters the Curie point. In this connection we have found in earlier work on N$_2$O$^{11}$ that the mean field model, on which our current analysis is based, is inaccurate at the highest temperatures, close to sublimation. We note therefore at the outset that the data for CO at 26 K are somewhat anomalous. This is presumably because these data are taken very close to the limiting temperature at which CO can be condensed$^{20}$ and, by implication, in a regime in which fluctuations in the system begin to dominate and the mean field model becomes increasingly inapplicable.

$\Delta\nu_{\text{B}}$ is estimated by plotting $\Delta\nu$ versus $1/T$ and extrapolating to zero, that is, indefinitely high temperature. This is an ad hoc approach, based first on simplicity and second on the high...
accuracy with which data points, excluding data for 26 K, lie on a straight line, which may therefore be readily extrapolated. $\Delta \nu$ versus $1/T$ is shown in Fig. 5 in which the slope is linear to within 0.6% and the associated intercept to within 0.3%, ignoring the 26 K data. This yields an intercept at 2.59 ± 0.01 cm$^{-1}$, which is the value assigned throughout to $\Delta \nu_0$. The resulting values of $\Delta \nu_S$ and $\xi = \Delta \nu_S/\Delta \nu$ are shown in columns 5 and 6 of Table 2.

### 3.2 A brief resume of the model for the spontelectric effect

At this stage we need to introduce the theoretical model which governs spontlectrics. A detailed description may be found in ref. 1. This model successfully describes the variation with deposition temperature, of the observed spontelectric field, in films of N$_2$O$_2$, N$_2$O, and CH$_3$Cl, CF$_2$Cl$_2$, and CFCI$_3$ films. The model is based on the concept that the net z-component of the electric field within a spontlectic film and normal to the plane of the film, $E_z$, is composed of two parts. The first is a local symmetrical part, defining the interactions which both bind layers to one another and dictate the molecular force field and thus molecular vibrational frequencies. The second is an asymmetrical part, due to the long-range field which permeates the film. The symmetrical part is expressed as a constant term plus a dipole–dipole term, proportional to $\langle (\mu_z)/\mu \rangle^2$, representing average intermolecular dipole–dipole interactions. Here, $\langle (\mu_z)/\mu \rangle$, the degree of dipole orientation, is the ratio of the average z-component of the dipole moment and the total dipole moment of the molecular species in the solid state, where the z-axis is perpendicular to the plane of the film. The $\langle (\mu_z)/\mu \rangle^2$ term reflects the fact that all dipole interactions, involving dipole-image charge, extended dipoles and arrays of dipoles, follow this squared relation.

We note that the symmetrical part of the contribution to $E_z$ is related to the ‘local field’ at any molecular site, as defined in standard texts.

The asymmetrical part of $E_z$ is described by $\langle E_{asym} \rangle (\mu_z)/\mu$ and is equal to the observed spontlectric field. This term is found only in the description of spontlectrics, with no direct counterpart for any other form of material. This asymmetrical part resembles the Weiss field in ferromagnetism, which is assumed to be proportional to the magnetisation. Here, degree of dipole orientation for magnetisation and read polarisation field for the Weiss field. We emphasise that the polarisation field, that is, the spontelectric field, is self-generated within the spontlectric material. The polarization field acts in opposition to the symmetrical part and represents the long-range field created by the average dipoles and experienced by an average dipole. Note that this description highlights the non-linearity of the interactions involved.

The spontlectric field is a result of the macroscopic polarization, $P$, of the film of CO. In this connection, note the absence of any free charges in the film. If we write $P$, as the polarization of a perfectly oriented system of dipoles, that is for $\langle (\mu_z)/\mu \rangle = 1$, then $P$ is given in the limit of point dipoles by $P = \mu / \Omega$, where $\mu$ is the dipole moment of CO in the solid state, and $\Omega$ is a parameter related to the effective molecular volume of the CO molecule. So the true polarization is given by $\langle (\mu_z)/\mu \rangle \mu / \Omega$. Now $P = \varepsilon_0 E_0$, where $E_0$ is the spontelectric field, noting that the use of $\varepsilon_0$ is appropriate since the dielectric effect of the medium has already been subsumed into the value of $\mu$ (see below). Thus $E_0 = \langle (\mu_z)/\mu \rangle \mu / \varepsilon_0 \Omega$ and $\langle E_{asym} \rangle = \mu / \varepsilon_0 \Omega$ or $4\pi \mu / \Omega$ in atomic units. $\Omega$ is treated subsequently as a parameter to be obtained through analysis of the experimental data, effectively replacing the parameter $\langle E_{asym} \rangle$.

Hence, using atomic units throughout,

$$E_z = \langle E_{asym} \rangle \left[ 1 + \frac{\langle (\mu_z)/\mu \rangle^2}{2} \right] - E_0$$

where $\langle E_{asym} \rangle$ and $\xi$ are taken to be temperature independent parameters. The temperature dependence of the spontlectric field enters in general through the temperature dependence of both $\langle E_{asym} \rangle$ and $\langle (\mu_z)/\mu \rangle$. The $\langle (\mu_z)/\mu \rangle^2$ term in eqn (1) may be interpreted as a measure of the tendency of one dipolar species to restrict the angular motion of another, a ‘locking’ term or, as it is sometimes called, a ‘frustration’ term.

Mean field theory gives an implicit expression for $\langle (\mu_z)/\mu \rangle$, yielding the familiar Langevin function for orientational interactions:

$$\frac{\langle (\mu_z)/\mu \rangle}{\mu} = \coth \left( \frac{E_0 \mu}{T} \right) - \left( \frac{E_0 \mu}{T} \right)^{-1}$$

where $T$ is the deposition temperature of the layer of material and the Boltzmann constant is unity in atomic units. The dipole moment of CO in the solid state is reduced from that in the gas phase through depolarization in the environment of other CO species according to:

$$\mu = \frac{\mu_0}{1 + 2k/s}$$

where $s$ is the average spacing between successive layers, $x$ is the molecular polarizability of CO (13.159 au), $k = 11.034^{46}$ and $\mu_0$ is the gas phase dipole moment of CO (= 0.122 D). $s$ is estimated from the diameter of isoelectronic N$_2$ to be 0.339 nm or 6.406 au.
3.3 An explicit expression for $\Delta \nu_B$

The purpose here is to derive an expression for the contribution of $\Delta \nu_B$ to the LO–TO splitting, in terms of the spontelectric parameters introduced in Section 3.2. The first task is to combine eqn (1) and (2) to obtain an explicit expression for the degree of dipole orientation, $\langle \mu_z \rangle / \mu$.

We approximate the coth function in eqn (2) by expansion to first order, writing that $\coth(x) \approx 1/x = 1/3x$. Using values relevant to the present context, this approximation is accurate to better than one part in $10^2$. We then obtain

$$\frac{\langle \mu_z \rangle}{\mu} = \frac{3T - 2\sqrt{2}\zeta \langle E_{\text{sym}} \rangle \langle E_k - \langle E_{\text{sym}} \rangle \rangle + 9T^2}{2\mu_0 \langle E_{\text{sym}} \rangle}$$

(4)

We now set out to express the LO–TO splitting in terms of spontelectric parameters. The internal electronic structure of the individual molecules, influenced by electrostatic effects from neighbouring molecules, is responsible for the force field associated with TO modes. Thus, $\langle E_{\text{sym}} \rangle [1 + \zeta (\langle \mu_z \rangle / \mu)^2]$, the first term of eqn (1), may be regarded as the average effective electric field at any molecule, giving rise to the force field which determines the value of $\nu_{\text{CO}}$ in the TO mode. However, the occurrence of a spontelectric field adds an additional potential in the direction normal to the plane of the film, shifting the LO vibrations to yet higher wavenumber than in intrinsic LO–TO splitting. Thus, the force field for LO includes an additional term involving the torque exerted on the molecular dipole in the spontelectric field. The field involved in this additional term has the form of the projection of the spontelectric field onto the direction in which the average dipole points, that is $E_0 (\langle \mu_z \rangle / \mu)$. Dipole orientation also influences the force field dictating the TO mode, via the term $\langle E_{\text{sym}} \rangle [\zeta (\langle \mu_z \rangle / \mu)^2]$ in eqn (1) and the coupling of $\langle \mu_z \rangle / \mu$ to $E_z$ in eqn (2).

$$\Delta \nu_S = -\frac{\mu_0 \nu_T E_0 T \left[ E_0 \left( 3T + 4\pi \mu^2 \zeta \langle E_{\text{sym}} \rangle (E_k - \langle E_{\text{sym}} \rangle) + 9T^2 \right)^{1/2} \right]}{2\langle E_{\text{sym}} \rangle \langle E_k - \langle E_{\text{sym}} \rangle \rangle + 9T^2}$$

(5)

onto the direction in which the average dipole points, that is $E_0 (\langle \mu_z \rangle / \mu)$. Dipole orientation also influences the force field dictating the TO mode, via the term $\langle E_{\text{sym}} \rangle [\zeta (\langle \mu_z \rangle / \mu)^2]$ in eqn (1) and the coupling of $\langle \mu_z \rangle / \mu$ to $E_z$ in eqn (2).

$$\Delta \nu_S = \frac{2\pi \mu^2 \zeta}{\Omega_0 \langle E_{\text{sym}} \rangle (4\pi \mu^2 + 3\Omega T)} \left\lbrace \begin{array}{l} 4\pi \mu^2 - \left[ 9\Omega^2 T^2 + 4\mu^2 \left( 4\pi \mu^2 - \Omega^2 \zeta \langle E_{\text{sym}} \rangle^2 + 6\pi \Omega T \right)^{1/2} \right] - 3\Omega T \end{array} \right\rbrace$$

(6)

Clearly the LO and TO modes possess two different effective force constants, reflecting the different force fields associated with LO and TO modes. Let $k$ be the force constant associated with a fictitious solid, in the absence of either the spontelectric effect or effects leading to the intrinsic LO–TO splitting. Then in a real solid, two force constants may be encountered, $k - \delta_T$ and $k + \delta_T$, where, introducing the harmonic approximation, $\nu_T \approx (k + \delta_T)^{1/2}$ which is equivalent to eqn (4). We note that since $E_0 = (\langle \mu_z \rangle / \mu) \mu / \nu_0 \Omega$, we can write

$$\frac{\langle \mu_z \rangle}{\mu} = \frac{\Omega E_0 \sqrt{4\pi \mu}}{2\mu_0 \langle E_{\text{sym}} \rangle}$$

(7)

which is exactly equivalent to eqn (7) but in which we have used $E_0 = (\langle \mu_z \rangle / \mu) \mu / \nu_0 \Omega$, inserting $\langle E_{\text{sym}} \rangle = 4\pi \mu / \nu_0 \Omega$ (Section 3.2) and we have employed

$$\langle \mu_z \rangle = \frac{\langle E_{\text{sym}} \rangle \mu + 3T \left[ 1 - \sqrt{(1 + \langle E_{\text{sym}} \rangle / \nu_0 \Omega)^2} - 4\mu^2 \zeta \langle E_{\text{sym}} \rangle^2 / 9T^2 \right]}{2\mu_0 \langle E_{\text{sym}} \rangle}$$

(8)

which is equivalent to eqn (4). We note that since $E_0 = (\langle \mu_z \rangle / \mu) \mu / \nu_0 \Omega$, we can write

$$\langle \mu_z \rangle / \mu = \Omega E_0 / 4\pi \mu$$

(9)
spontenelectric data\cite{1,7,8} that of the terms set out in Table 1, $\mu, \zeta$, $\langle E_{sym}\rangle$ are constant with temperature of deposition. In contrast to earlier work, we do not assume that $\langle E_{asym}\rangle$, or equally $\Omega$, is constant with temperature of deposition. In Section 3.4.2, it is found that $\Omega$ may vary with temperature.

We proceed through four distinct steps.

(i) The first step is to make an estimate of $\langle E_{sym}\rangle$. We use eqn (8) to write down pairs of simultaneous equations, where each pair refers to a specific combination of two temperatures. These are then solved to yield six values of $\langle E_{sym}\rangle$ and $\Omega$. The average value of $\langle E_{sym}\rangle$ is used in subsequent calculations.

(ii) In the second step, we use eqn (8) at all five different experimental temperatures and solve for $\Omega$, using the average value of $\langle E_{sym}\rangle$ from step (i).

(iii) We solve eqn (7) to yield values of the spontenelectric field, $E_{sym}$.

(iv) We use eqn (10) to determine values of $\langle \mu_z\rangle/\mu$.

### 3.4 Extracting values of the spontenelectric field as a function of deposition temperature

We now seek to use eqn (8) to obtain parameters of the spontenelectric field. In order to simplify this procedure, we first take note that the value of $\Delta \zeta$ in eqn (8) is insensitive to the value of $\zeta$. We know from earlier work that $\zeta$ can assume a wide range of values,\cite{1} from a few tens to $>10000$. However, numerical tests on eqn (8) show that variation of $\zeta$ over this range causes $\langle E_{sym}\rangle$ to vary by <5%, given the values of the other parameters encountered here. This suggests the simplifying measure, which we adopt here, that $\zeta$ assumes some chosen value in the expected range and that the quantitative results of subsequent analysis are essentially independent of the value of this choice.

Physically, $\zeta$ determines the strength of the dipole locking term, as we have mentioned. Since the dipole moments in the solid state of CO and N$_2$O are essentially the same, respectively 0.0785 D and 0.0786 D, and the layer spacing is again similar, with $S_{N_2D} \sim 6.05$ au (0.32 nm) and $S_{CO} \sim 6.4$ au (0.34 nm), we have chosen to use the same value of $\zeta$ as derived from experimental data in ref. 1. Thus we adopt $\zeta = 43.8$, recognizing that the value could be reduced or increased by a factor of ten or more without making any significant difference to our conclusions.

#### 3.4.1 Results of step (i): a solution for $\langle E_{sym}\rangle$.

A typical example of a pair of simultaneous equations in $\langle E_{sym}\rangle$ and $\Omega$, derived from eqn (8) for $T = 20$ and 24 K, are as follows:

\begin{equation}
1.44 = \frac{26.77 + \Omega \left[ 0.423 - 2228 \sqrt{3.61 \times 10^{-8} + (1.44 \times 10^{-4} + 4.57 \times 10^{-6})/\Omega^2 - 1.34\langle E_{sym}\rangle^2} \right]}{\Omega\langle E_{sym}\rangle(\Omega + 52.7)}
\end{equation}

and

\begin{equation}
1.20 = \frac{19.76 + \Omega \left[ 0.375 - 1644 \sqrt{5.19 \times 10^{-8} + (1.44 \times 10^{-4} + 5.48 \times 10^{-6})/\Omega^2 - 1.34\langle E_{sym}\rangle^2} \right]}{\Omega\langle E_{sym}\rangle(\Omega + 52.7)}
\end{equation}
molecular volume. A recent paper\textsuperscript{27} reported the correlation between polarizability, $\alpha$, and molecular volume and suggested the empirical relationship $\alpha = 0.0086161 \times V^0.33$. Using $\alpha = 13.159$ au, we find $\Omega = 244 \pm 30$ au. The average of $\Omega_3$ in Table 3 is 273, favouring the choice of parameters associated with $\Omega_2$. At the same time, one should beware of equating $\Omega_2$ rigidly with the molecular volume, since results in Table 3 show that the value of this parameter changes rapidly with temperature, given the assumptions made in our analysis. We also note that $\Omega_1 \sim 5$ to $6 \times 10^4$ yields unphysical degrees of dipole orientation, greater than unity.

### 3.4.3 Results of step (iii): solutions for the spontlectric field

Solutions of eqn (7) for the spontelectric field are shown in Table 4, using data from Table 2 and again using the average value of $\langle \Omega_{\text{sym}} \rangle = 4.48 \pm 0.21 \times 10^7$ V m$^{-1}$ from step (i). Results in Table 4 show two values for the spontlectric field for each deposition temperature, the upper branch value $E_{\ell u}$ and the lower, $E_{\ell u}$.

### 3.4.4 Results of step (iv): solutions for the degree of dipole orientation

Substituting values of $E_{\ell u}$, $E_{\ell u}$ and $\Omega_2$ into eqn (10) yields values of $\langle \mu_z \rangle / \mu$ for each temperature. Results, subscripted $u$ and $l$ are shown in Table 5, for the upper and lower spontlectric branches respectively.

In Table 5, consider first the upper branch, $E_{\ell u}$. This is well-behaved, in the sense that the spontlectric field declines as the temperature of deposition is raised and this is accompanied by a decline in dipole orientation. The lower branch, $E_{\ell u}$, by contrast shows the peculiar feature that the spontlectric field increases with temperature, accompanied however by the expected fall in dipole orientation. It may be possible to pursue the grounds for this behaviour through differentiation of eqn (7) with respect to $T$. This is not developed further here.

We now briefly compare our present results with data for the vibrational Stark shift of CO in the gas phase.\textsuperscript{28} Following results reported in ref. 28, the fields presented in Table 5 would create a Stark shift of $0.93 \pm 0.04$ cm$^{-1}$ and $0.19 \pm 0.08$ cm$^{-1}$ for the upper and lower branch electric fields respectively, in the gas phase. We observe a shift of $\sim 0.67$ cm$^{-1}$, taking the average value of $\Delta \nu/2$. This serves to illustrate that our derived values of field are consistent with gas phase Stark data. However it leaves open the question of which of the two branches is relevant to the present experiment. On the basis that the Stark shift is likely to be lower in the solid state than in the gas phase due to depolarization, the above estimates would favour the higher value of spontlectric field, associated with the upper branch in Table 6. A comparison with data for N$_2$O, in Section 4 below, supports this suggestion.

We note that the double-valued, or multi-valued, behaviour, seen here in spontelectrics, is characteristic of non-linear, non-local systems, such as photonic crystals\textsuperscript{29} where the bistability arises through non-linear Maxwell equations, to mention one of very many and varied examples. Moreover, this bimodal behaviour seen in spontelectrics – given that it is not some mathematical artefact – suggests new avenues of research. For example, the presence of two states represents a bi-metastable system, which may allow switching through an applied electric field. For the present, however, we do not explore this behaviour further, leaving it to future more detailed investigation. From hereon, we consider only the upper branch values of the spontlectric effect in Table 5.

### 4. Comparison with other spontlectric materials

The value of the spontlectric field in the upper branch of CO (Table 6) is comparable to that encountered in N$_2$O, the latter around a deposition temperature of $\sim 60$ K. The values of the...
degree of dipole orientation in the upper branch are also similar to those encountered in N₂O in films, deposited between 52 and 65 K, which possess (μₐ)/μ = 0.0665 and 0.0204 respectively. The temperature of 65 K is close to the highest temperature at which deposition is possible for N₂O, much as is 26 K for CO. These remarks suggest some inherent similarity between the two species in their spontelectric properties. This may derive from their very similar values of dipole moment in the solid state, mentioned earlier.

The value of the spontelectric field for CO, expressed as mV of surface potential added per added monolayer, may be compared with values known for additional species, although typically at 40 K rather than 20 K. This comparison is illustrated in Table 6, where values for other species are taken from ref. 1. Values of the degree of dipole alignment are also shown in Table 6, where these are known. The significance of the data in Table 6 is to show that the characteristics, deduced for CO from RAIRS, fall into a bracket already encountered for spontelectrics. It is interesting to note that propane, in Table 6, explicitly shows a double valued nature, switching from an upper branch to a lower branch at a film thickness of ~2500 ML.

In addition, we can check our present method using data for N₂O. In ref. 8, as described in the introduction, we calculated the LO–TO splitting using spontelectric parameters for N₂O, known from earlier work based on direct measurement of the surface potential using the electron beam method. We then compared the measured LO–TO splitting at a variety of deposition temperatures, showing that the spontelectric Stark effect reproduces the LO–TO splitting data. We now test the use of the measured LO–TO splitting in N₂O, recorded in ref. 8, in reproducing the known degree of dipole orientation and the corresponding spontelectric field at temperatures between 48 and 60 K, using the analysis developed here.

We use the following parameters for N₂O: ζ = 75, Ω = 255 au, μ = 0.0785 D, the latter corresponding to a layer separation of 0.32 nm, taken from ref. 1. The value of Ω is that used in both ref. 1 and 8 and has been adopted here since we seek a comparison between different methods on the same footing. With data for νᵣₑ, ζ and Δνₛₒₜ for N₂O taken from RAIRS data in ref. 8, we find, rearranging eqn (7), that the average value of ⟨Eₜᵢₑ₅⟩ is 4.69 ± 0.19 × 10⁴ V m⁻¹. This agrees within experimental error with the value of 4.57 ± 10⁴ V m⁻¹, reported in ref. 8. Note that the latter value was derived from fitting experimental data of surface potentials vs. temperature of deposition between 48 and 60 K for N₂O, reported in ref. 1, using the model outlined in Section 3.2.

We now use eqn (9) to calculate (μₛ)/μ as a function of temperature. Here Eₜ = ⟨Eₜᵢₑ₅⟩/μ = (μₛ)/μ/ζΩ, from eqn (10), yields the spontelectric field. Values are shown in column 5 of Table 7. The temperatures of deposition in Table 7 are selected to cover the range of values of experimental data available, using both the electron beam method and RAIRS. The second column of Table 7 shows values of (μₛ)/μ estimated from the present model, based, that is, upon RAIRS measurements. The third column shows results obtained using spontelectric parameters obtained by fitting to surface potential data, using the theory outlined in Section 3.2. The fourth column presents values of dipole orientation derived from observational data in ref. 4. The three succeeding columns give corresponding values of the spontelectric field. The most telling comparison, for validation of the present model, is of values in columns two with those in three and those in five with those in six. Agreement with observational data, in columns 3 and 7, deteriorates at higher deposition temperature, reflecting the less accurate fitting for (μₛ)/μ in this regime.

Results in Table 7 validate the present method as it is applied here to CO. We note that values of the spontelectric field for N₂O in Table 7 appear consistently 5 to 7% higher, using the RAIRS data, than through numerical fitting to surface potential data.

### 5. A meeting point for condensed matter science and astrophysics

One may ask whether the existence of spontelectric CO has any consequence in the natural Universe. The reply is that this may well be the case, since interstellar dust grains could acquire a substantial polarization charge on the grain surface. This possibility is briefly explored below, for the most part qualitatively.

Without elaborating at length, it is well-established that interstellar grains in cold pre-stellar cloud cores are coated with an outer layer of quite pure solid CO (e.g. in Taurus). The extent of this layer may be 50 to 60 ML, for example in the molecular core in L1544 in Taurus. We suppose below that the low rate of accumulation of CO in space has no influence on the spontelectric nature of the film formed.

Using figures from the present work, the spontelectric effect in CO at 20 K would generate a polarization potential of 0.7 V on such a grain, given a thickness of 50 to 60 ML. This potential is equivalent to a polarization charge of 50 electronic charges, noting that a grain is generally regarded as possessing typically one negative charge on its surface. Following ref. 31, which suggests that the positive O-end of CO would protrude from the surface, the polarization charge would be 50 positive.

It is evident that such a polarization charge would tend to attract electrons to the surface, altering the gas phase abundance of electrons. Using a density of grain material of 2 g cm⁻³ and a grain abundance by mass of 1.3%, the proportion of electrons removed to the surface of the grains may be shown to be given by 1.06 × 10⁻²⁵ nₛₐ₅ / (a²), where nₛₐ₅ is the number of...
monolayers of CO, \( z \) is the degree of ionization and \( a \) is the grain radius. Values of \( z \) range\(^{32} \) over a few \( 10^{-9} \) to \( 10^{-8} \). Thus, say, if \( a = 0.025 \) \( \mu m \) and \( z = 10^{-8} \), then the degree of depletion of gas phase electrons would be \( > 80\% \). This, in turn, would influence the abundance of chemical species in the interstellar medium. Molecular species, in particular water, are important for their cooling effect in a cloud under gravitational collapse and thus in maintaining such collapse approximately isothermal. There is an additional link to the chemistry through the nature of magnetic shocks, whose characteristics are strongly dependent on the gas phase abundance of electrons: see ref. 33 and 34 and references therein. Such shocks form both in star forming regions and close to the surface of nascent protostars.

Whilst the above suggests that the inclusion of spontelectric grains may have an influence on the chemistry and physics of the gravitational collapse of protostellar cores, a proper analysis requires a more complex – and interesting – kinetic approach especially isothermal. There is an additional link to the chemistry of a magnetic field, \( B \) in L1544, which influences the ambipolar diffusion. Coupled further with this, are the effects of electron adsorption will increase as the surface becomes more polarized.\(^{35} \) At the same time, electro-neutrality of the plasma drag will pull positive ions towards the grains, leading to ambipolar diffusion. Coupled together with this, electron drag will pull positive ions towards the grains, leading to ambipolar diffusion. Coupled further with this, are the effects of a magnetic field, \( \sim 0.5 \) mG, which influences the motion of electrons. Future work in astrophysical modelling of star-forming clouds should therefore involve a time-dependent coupled magnetohydrodynamic and kinetic model in order to predict the time dependent effects of the presence of polarized grains in astrophysical plasma.

6. Concluding remarks

Our conclusions may be summarised as follows:

(i) As a standalone technique, RAIRS can be used to establish the spontelectric character of films. All that is required is the observation that the LO–TO splitting shows a measurable dependence on the film deposition temperature, and possibly also upon annealing when the effect is large enough. Data at a variety of deposition temperatures should then yield the strength of the spontelectric field and the degree of dipole orientation, as a function of film temperature. This would allow an exploration of the range of species which are spontelectric. RAIRS does not however inform us in which direction the field is oriented.

(ii) Our results call for new developments in modelling the chemistry and physics of pre- and protostellar cores, as part of the current extensive research into the nature of low-mass star formation and of planet-forming disks around young stars.\(^{36} \)

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